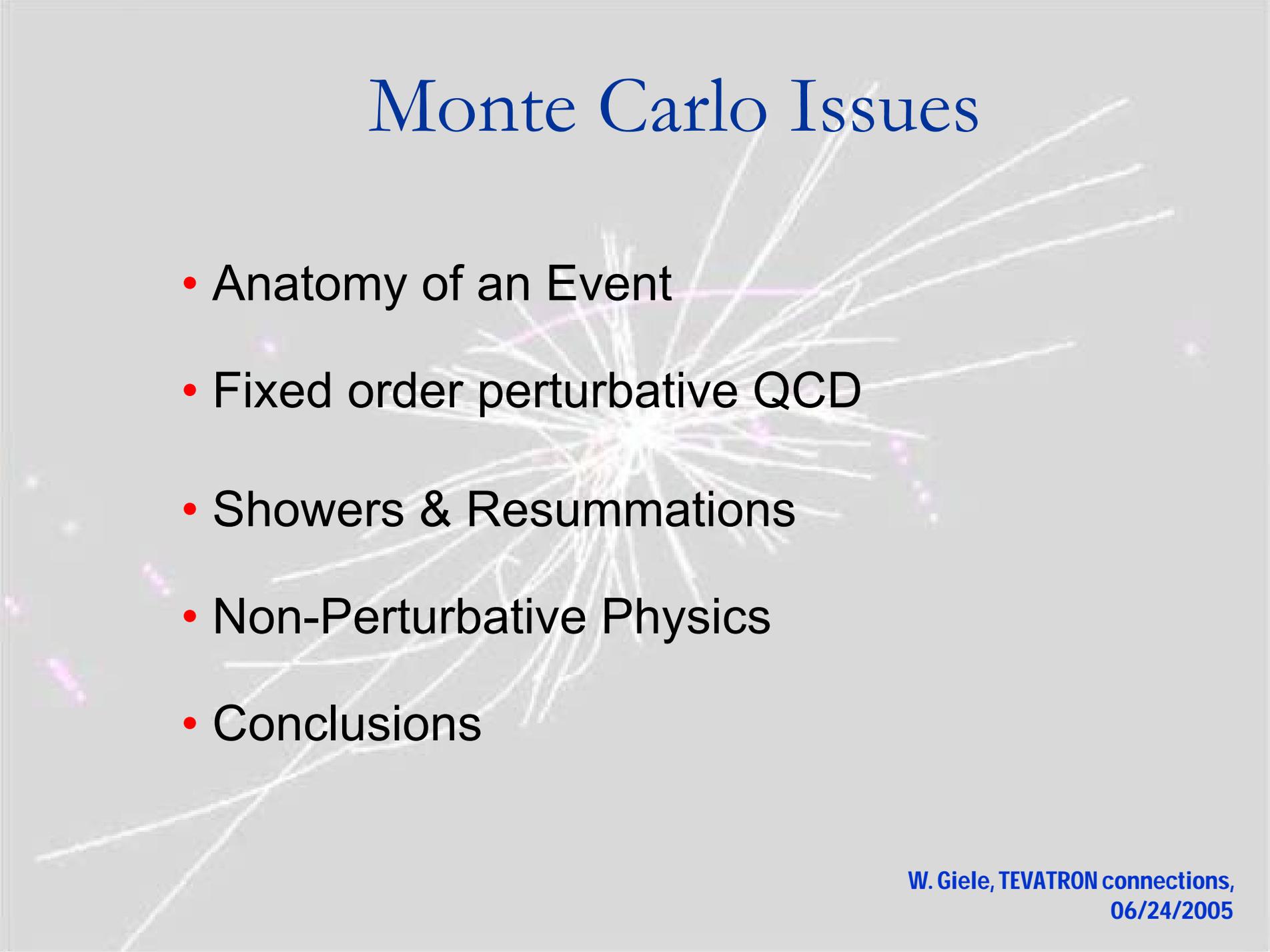


# Monte Carlo Issues

A background image showing a complex network of white and purple lines radiating from a central point, resembling particle tracks or a starburst pattern.

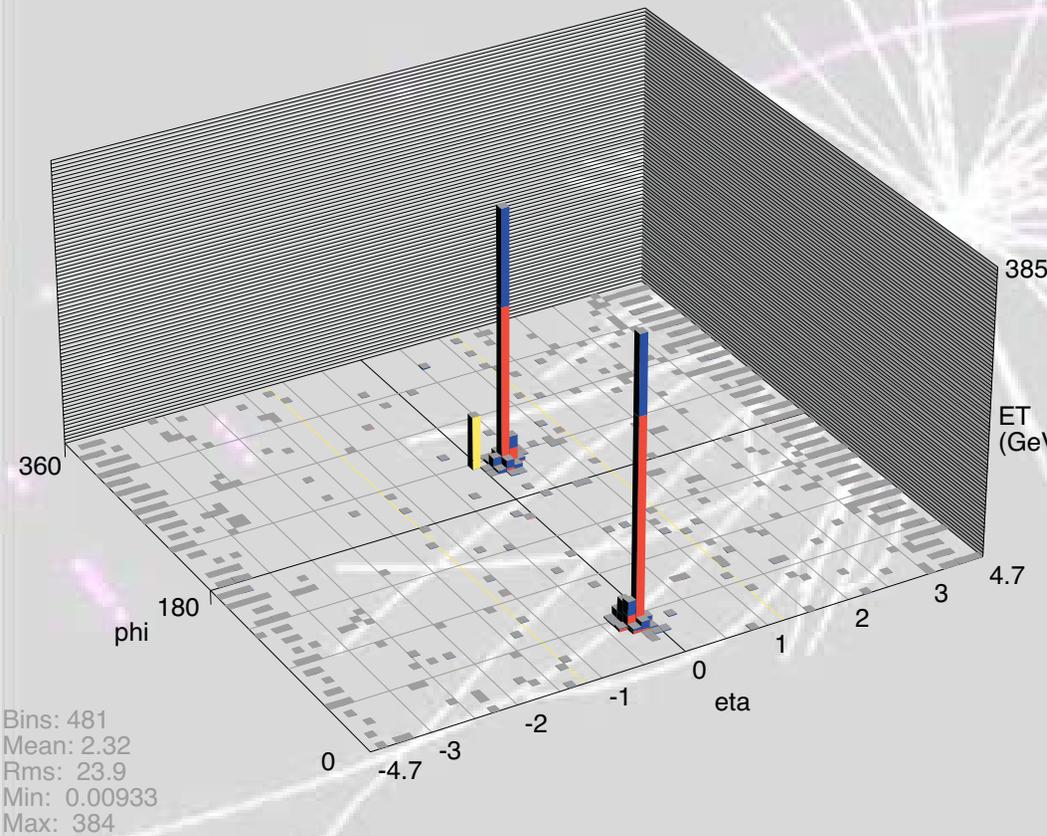
- Anatomy of an Event
- Fixed order perturbative QCD
- Showers & Resummations
- Non-Perturbative Physics
- Conclusions

# Anatomy of a (collider) Event

- The experiment is a series of events defined by a resolution scale (imposed by analysis such as jet resolution or detector resolution).
- The MC generator is a series of events also defined by a resolution scale (cluster resolution or hadronization scale).
- For theory and experiment to agree both series of events should be statistically equivalent.

# Anatomy of an Event

Run 178796 Event 67972991 Fri Feb 27 08:34:03 2004



mE\_t: 72.1  
phi\_t: 223 deg

- For a high momentum transfer event the event looks like a  $2 \rightarrow 2$  scattering event up to reasonably small resolution scales.
- A simple leading order description ( $gg \rightarrow gg$ ,  $gq \rightarrow gq$ , ...) will do a good job simulating the di-jet correlations.
- Pushing towards a smaller and smaller resolution scale will reveal more structure (in the MC we will get “large logs” of the resolution scale).
- Next-to-leading order will describe the additional structure (average description)
- Reducing the resolution scale even further will require even higher orders and eventually we will need resummations (i.e. shower MC's)

# Anatomy of an Event

- At leading order the parton represents the average behavior of the hadronic clusters, provided the resolution scale is sufficiently large to encapsulate the cluster.

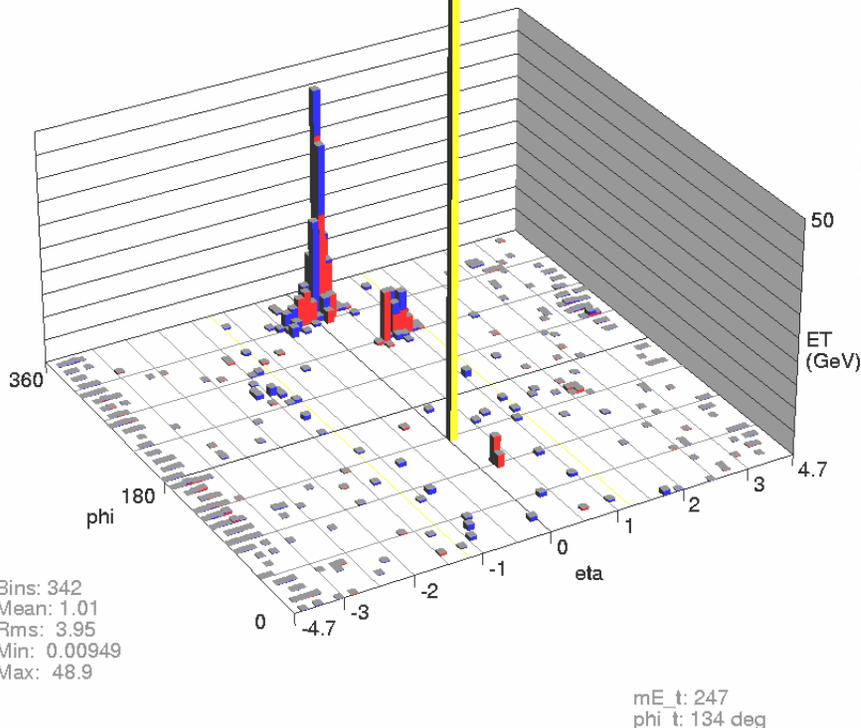
- Leading order cannot describe well the absolute probability (normalization) but will work reasonably well for event shapes (relative probability)

- Higher orders will be able to predict normalizations and depend on cluster shapes (i.e. defines its own resolution scale).

- To go to a calorimeter level scale resolution one would need resummed partonic showers.

- To go beyond hadronization scale one would need a hadronization model. Detector response will depend on hadronic correlations. This requires the MC to model the event all the way down to hadrons.

Run 180952 Event 51963432 Tue Mar 16 18:07:08 2004



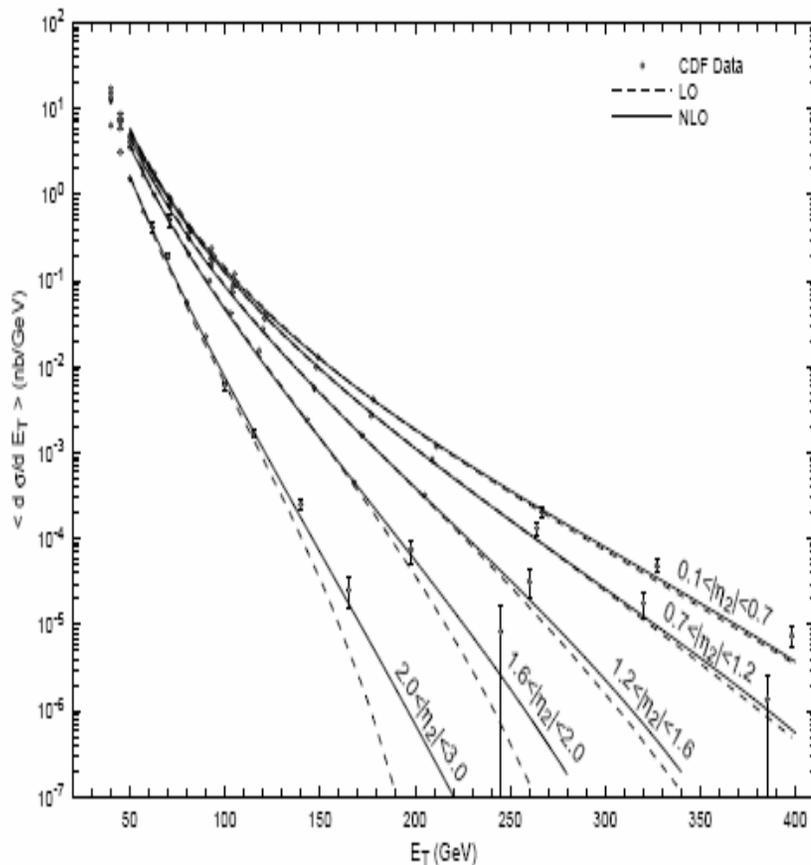
# Anatomy of an Event

- The MC event goes through 3 steps:
  - Large resolution scale  $Q$ , such that  $\alpha_s(Q) \ll 1$   
This results in a fixed order perturbative expansion (deterministic)
  - Small resolution scale  $Q$  such that  $\alpha_s(Q) \times \text{Log}(Q) \approx 1$   
This results in a resummed parton shower MC (some arbitrariness, but controllable)
  - Hadronization scale, controlled by non-perturbative physics. Only phenomenological models exist. Needs data to fit to...
  - What are the issues in a chain of tools like **MCFM**  $\rightarrow$  **MC@NLO**  $\rightarrow$  **HERWIG** ?
  - How do we improve on this set of tools...

# pQCD

- This is the first step in any MC generator.
- Given the Feynman rules the answer to a given scattering process is unique.
- At leading order (“tree graphs”) we can calculate all necessary scattering amplitudes.  
(started with VECBOS and evolved over the years to products like ALPHGEN, MADGRAF,...)

# pQCD



- LO should describe the correlations between the average directions of the hadronic clusters well (provided appropriate large resolution)
- Difference between LO and NLO is small (except for normalization)
- NNLO and beyond, including partonic shower will/should not change anything
- Unfortunately the calorimetric detector response depends on the hadronic content of the clusters
- This makes the data comparison as shown here sensitive to hadronic effects and its interactions with the detector.
- Consequently the ability to add shower MC to fixed order is important.

# pQCD

- Why is NLO (one-loop) so important ?
  - We become sensitive to the meaning of the resolution scale, i.e. the calculation becomes sensitive to the choice of scale.
  - It gives the first reliable estimate of the normalization of the cross section
  - First sensitivity to the notion of dipoles (instead of the naïve jet). I.e. a hadron is part of a dipole, not a jet → new class of “jet algorithms” (which will be required as precision of measurement increases)

# pQCD

- Progress in NLO calculations has been slow
  - All **2 to 2** NLO processes (including quark masses)
  - Some **2 to 3** NLO processes (no quark masses)
  - No **2 to 4** NLO processes
- Analytic calculations are running into complexity problems and progress daunting
- Alternatives seems to be needed giving a systematic calculational procedure:  
The Samper project (c++, f95, f77)  
(**S**emi-numerical **AM**plitude **E**valuato**R**)

# pQCD

- Development for semi-numerical evaluation of one-loop calculations.
  - Detailed algorithmic method has been developed.
  - Program has been checked and is ready for **2 to 3** processes (no internal masses yet)
  - Will extend MCFM to:
    - **Di-boson + 1 jet** production
    - **Tri-boson production + 0 jet** production
    - **H + 2 jets** (with effective **Hgg** coupling)

# pQCD

- Currently implementing **H + 2 jets** production (K. Ellis, W. Giele, G. Zanderighi)
- Excellent numerical stability
- We were able to get an analytic result for **H + 4 quarks**, not for the other two processes.
- Example: **H + qq̄b + rrb̄**:
  - analytic:  
$$-46.7813035247351/e^2+(111.948110122775+18.3709749348328 i)/e+(120.012242523826-335.917283834563 i)$$
  - numerical:  
$$-46.7813035247350/e^2+(111.948110122775+18.3709749348302 i)/e+(120.012242523817-335.917283834578 i)$$
- Other processes implemented numerically, checked on gauge invariance etc. (~1 part in  $10^{11}$ )

# pQCD

- Next steps (programmatically approach):
  - Implementing internal masses. This will give:
    - $Q \bar{Q} + \text{jet}$
    - $Q \bar{Q} + V$
  - 2 to 4 processes. This will give access to a huge range of processes:
    - $Q \bar{Q} + Q \bar{Q}$  (eg top-pair + bottom-pair)
    - 4 jets
    - 3 jets +  $V$  (including mass effects for quarks)
    - 2 jets +  $V V$  (including mass effects for quarks)
    - .....

# Parton Showers

- The parton showers evolve from the resolutions of well separated clusters down to the hadronization scale.
- It is highly desirable to interface the pQCD calculations with the shower MC
- All current efforts are based on trying to interface with PHYTIA/HERWIG/...
- These shower MC in themselves correctly give the “leading log” behavior (i.e. resumming the most divergent resolutions scale logarithms)
- One runs in some issues when interfacing with pQCD calculations.

# Parton Showers

- Suppose we know the unresolvable factorization behavior of the pQCD matrix elements:  
$$A_n \simeq T_{n-1 \rightarrow n} A_{n-1}$$
- This soft/collinear factor contains the large logarithms we need to resum.
- How do we get a shower description without double counting?
- Suppose we know  $A_2, A_3, \dots, A_n$ , then

$$\sum_{m=2} A_m = A_2 + A_3 + \dots + A_{n-1} + S_n A_n$$

$$S_n = 1 + T_{n \rightarrow n+1} + T_{n \rightarrow n+1} T_{n+1 \rightarrow n+2} + \dots$$

# Parton Shower

$$\sum_{m=2} A_m = A_2 + A_3 + \dots + A_{n-1} + S_n A_n$$

$$= A_2 + A_3 + \dots + S_{n-1} A_{n-1} + S_n (A_n - T_{n-1 \rightarrow n} A_{n-1})$$

$$= A_2 + A_3 + \dots + S_{n-1} A_{n-1} + S_n \tilde{A}_n$$

$$= S_2 A_2 + S_3 \tilde{A}_3 + \dots + S_n \tilde{A}_n$$

with

$$S_m = 1 + T_{m \rightarrow m+1} S_{m+1}$$

$$\tilde{A}_m = A_m - T_{m-1 \rightarrow m} A_{m-1}$$

Note that the subtracted matrix element goes to zero as the resolution scale goes to zero on an event by event basis

# Parton Shower

- This is all we need...
- We shower off the subtracted matrix elements..
- No double counting...
- The subtracted matrix elements are finite at any resolution scale (over the whole phase space)..
- The subtraction function is the same as the exponentiated function  $\rightarrow$  strongly correlated
- The subtraction functions are well known

# Parton Shower

- However, this is not the HERWIG/PYTHIA type shower:
  - The subtracted matrix element only goes to zero in the soft limit averaged over many events...
  - This is not really a problem as far as the shower MC go on themselves (collinear correct event by event)
  - However this is a problem for interfacing with multiple pQCD matrix elements (need to modify the subtraction function for matrix element, while leaving exponentiated function unchanged)

# Parton Shower

- The “correct” shower is in fact based on  $2 \rightarrow 3$  branching (dipole branching) instead of the splitting ( $1 \rightarrow 2$  branching)
- An “engineering project” is underway to construct the desired shower MC: Higgs  $\rightarrow$  gluons (Giele, Kosower, Skands)
- After completion this will be build up to a full shower MC (VIRCOL shower MC)

# Hadronization Model

- Here the resolution scale is pushed below 1 GeV and individual hadrons are resolved
- This is a subject without much theoretical guidance
- To make systematic progress we need to understand the uncertainties in the pQCD and parton shower part well
- This is still in the future, for now both PYTHIA and HERWIG have QCD inspired phenomenological models

# Conclusions

- We are busy addressing current issues:
  - Samper project: pushing the NLO calculations to  $2 \rightarrow 4$  processes and beyond
  - VIRCOL project: improved parton shower MC with exact matching to LO/NLO/... matrix element
- Within the time span of run II we could complete:
  - All  $2 \rightarrow 3$  processes at NLO (including quark masses... e.g.  $VVV$ ,  $VV+\text{jet}$ ,  $V+QQ\text{bar}$ ,  $\text{jet}+QQ\text{bar}$ )
  - A VIRCOL shower which incorporates LO matrix elements (VIRTEV)