

Measurement of the Longitudinal Proton Structure Function in Diffraction

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HERA Accelerator

- the one and only ep accelerator in the world (since 1992)
- data taking ended on 02.07.2007
- recently, there is no other ep collider to be built



$$E_p = 920 \text{ GeV}$$

$$E_e = 27.5 \text{ GeV}$$

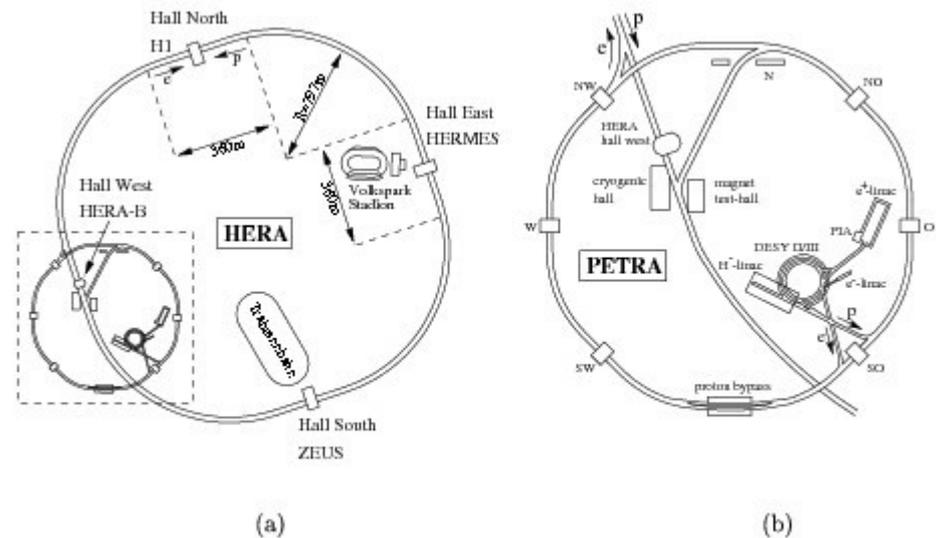
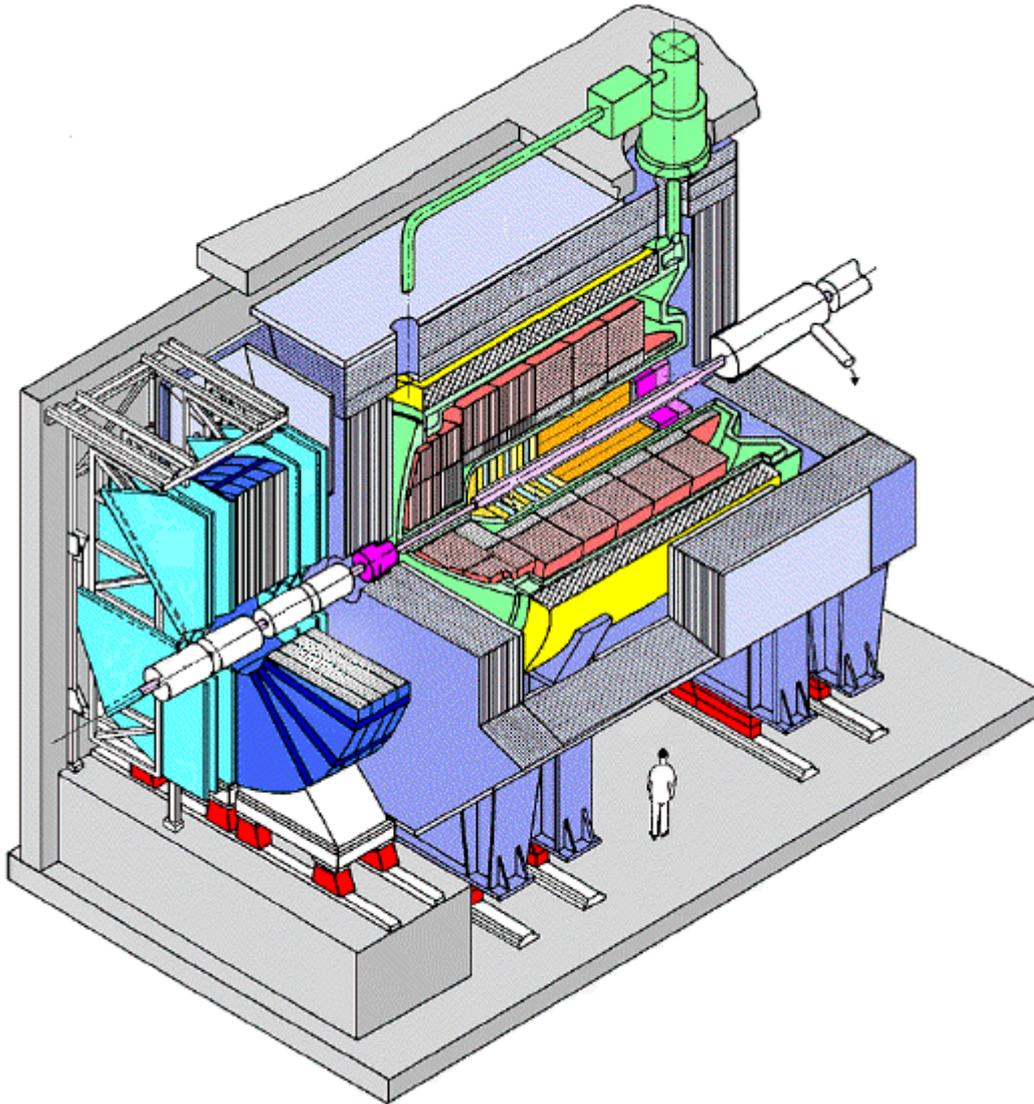


Fig. 3.1. HERA (a) and its pre-accelerators (b).

H1 Experiment



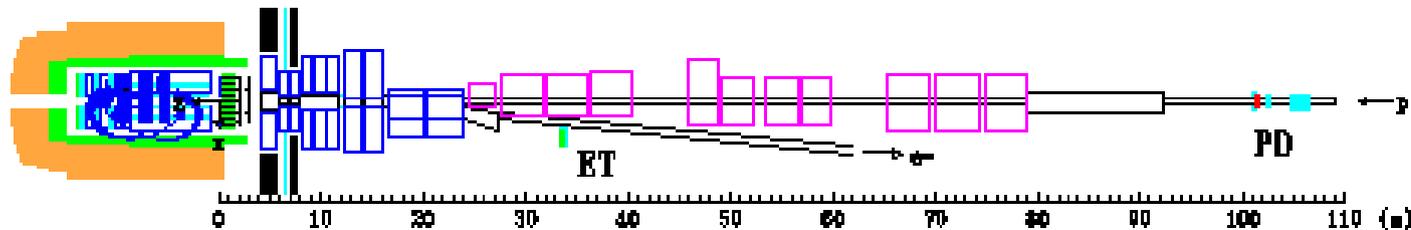
H1 Collaboration at DESY, Notkestr.85, D-22607 Hamburg, Germany



- tracking detectors
 - silicon detectors
 - drift chambers
- calorimeters
 - LAr, spaghetti calorimeter
- muon chambers
- forward detectors
 - FMD, PRT
 - FPS, VFPS
- and others

Cross Section Measurement

- number of events is proportional to luminosity $N = \sigma L$
- luminosity can be calculated out of the beam properties
- but it is more likely to be measured by some well known process
- Bethe-Heitler process $ep \rightarrow e p \gamma$
 - calculable within QED
 - accuracy $\sim 1\%$

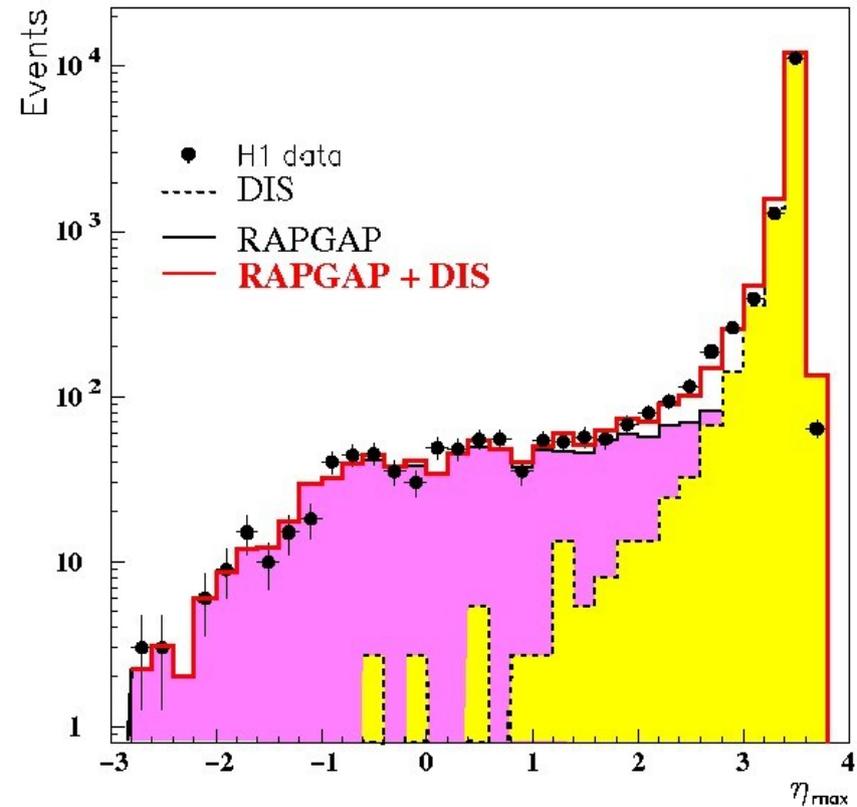
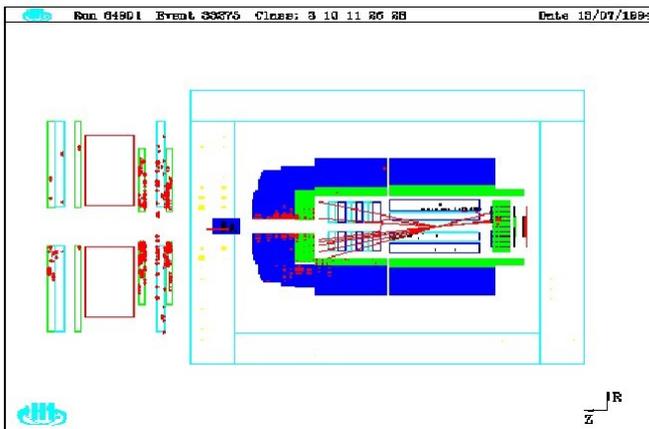
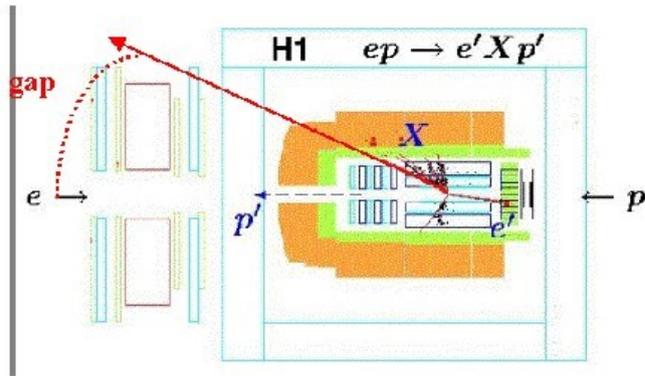


Rapidity Gap

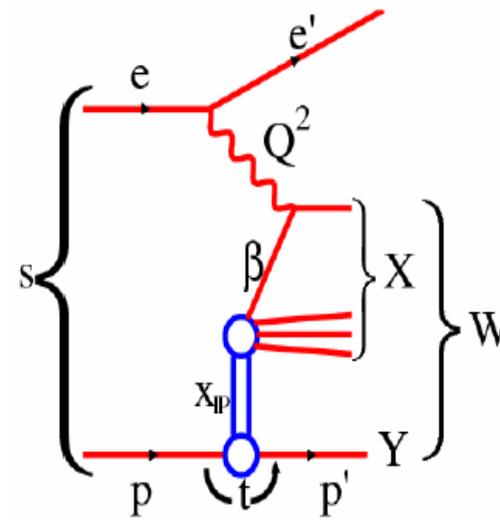
- rapidity
- pseudorapidity

$$y = \frac{1}{2} \log \frac{E + p_z}{E - p_z} = \frac{1}{2} \log \frac{(E + p_z)^2}{m^2 + p_t^2}$$

$$\eta = y|_{m=0} = \log \frac{E + p_z}{p_t} = -\log \left(\tan \frac{\theta}{2} \right)$$



Diffraction



~ 10% of DIS events
have a rapidity gap

$$t = (p - p')^2$$

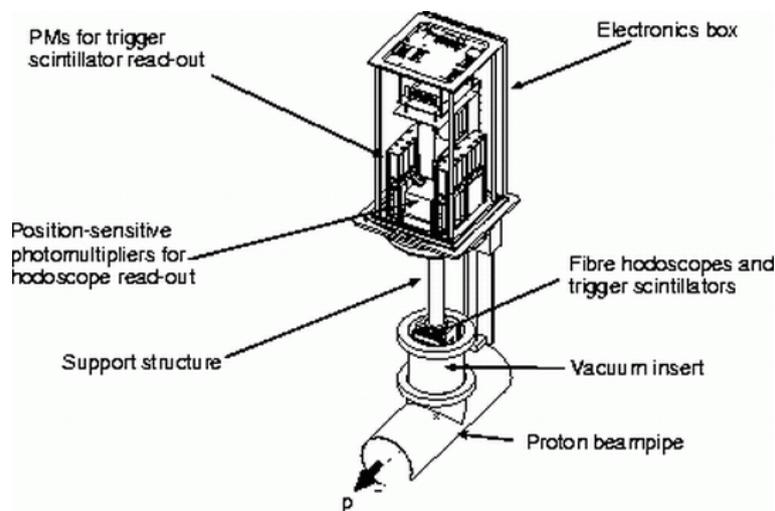
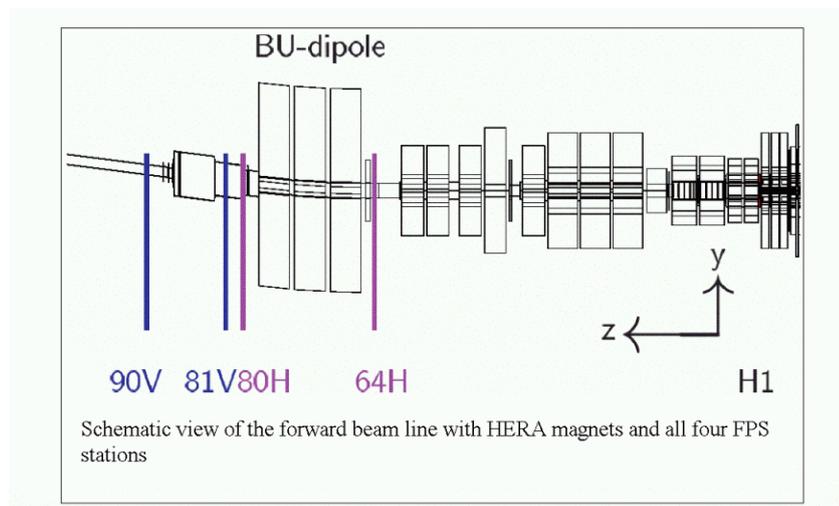
$$\beta = x_{quark/IP}$$

$$x_{IP} = x_{IP/proton}$$

- inelastic process where the proton remains intact
 $ep \rightarrow epX$
- rapidity gap reflects the exchange of colourless object with quantum numbers of vacuum (Pomeron)
- there is no colour flow between the final state X and the scattered proton Y

Forward Proton Detectors (FPS, VFPS)

- Roman pots detect scattered protons and measure their momenta
- an alternative to large rapidity gap cut to select diffractive data is to require a scattered proton in the Roman pots
- small acceptance makes the rapidity gap cut more favourable (for many analyses) as it is good to have larger statistics
- still, the Roman pots can be used to understand background better



Cross Section – Transversal and Longitudinal Component

$$\sigma = c \varepsilon^\mu \varepsilon^\nu W_{\mu\nu}$$

$$W_{\mu\nu} = -W_1(p, q) g_{\mu\nu} + W_2(p, q) \frac{p_\mu p_\nu}{M^2}$$

$$\varepsilon_L = \frac{1}{\sqrt{Q^2}} (q_3, 0, 0, q_0)$$

$$\varepsilon_L^2 = 1$$

$$q_3^2 - q_0^2 = Q^2$$

$$q_3^2 = v^2 + Q^2$$

$$\varepsilon_T = \frac{1}{\sqrt{2}} (0, 1, 0, 0)$$

$$\varepsilon_T = \frac{1}{\sqrt{2}} (0, 0, 1, 0)$$

$$\varepsilon_T^2 = -1$$

$$p = (M, 0, 0, 0)$$

$$\sigma = c \left(-W_1 \varepsilon^2 + W_2 \frac{(\varepsilon \cdot p)^2}{M^2} \right)$$

$$\sigma_T = c W_1$$

$$\sigma_L = c \left(-W_1 + W_2 \frac{q_3^2}{Q^2} \right) = c \left(-W_1 + W_2 \frac{Q^2 + v^2}{Q^2} \right) \rightarrow_{Q^2 \ll v^2} c \left(-W_1 + W_2 \frac{v^2}{Q^2} \right)$$

Measurement of the Longitudinal Proton Structure Function in DIS Using Two Different Centre-of-Mass Energies

$$\frac{d^2 \sigma}{dx dQ^2} = \frac{4 \pi \alpha^2}{Q^4 x} [F_2(1-y) + F_1 y^2 x]$$

$$F_2 = \nu W_2$$

$$F_1 = M W_1$$

$$R = \frac{F_2}{2xF_1} - 1 = \frac{W_2 \nu^2}{W_1 Q^2} - 1 = \frac{\sigma_L}{\sigma_T}$$

$$\frac{d^2 \sigma}{dx dQ^2} = \frac{4 \pi \alpha^2 Y_+}{Q^4 x} [F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2)]$$

$$Y_+ = 1 + (1-y)^2$$

$$\frac{d^2 \sigma}{dx dQ^2} = \Gamma(E, x, Q^2) [\sigma_T + \sigma_L \varepsilon(E, x, Q^2)]$$

- x, Q^2 fixed

- y vary $Q^2 = s x y$

- proton structure function F_2 is measured according to this formula
- structure functions F_2 and F_L cannot be separated unless you vary y
- without varying y the measurement of neither of the structure functions is not direct, it is based on some assumptions on one of them

Measurement of the Longitudinal Proton Structure Function in Diffraction

$$\frac{d^3 \sigma^{ep \rightarrow eXY}}{dx_{IP} d\beta dQ^2} = \frac{2\pi\alpha^2}{\beta Q^4} Y_+ \sigma_r^D(x_{IP}, \beta, Q^2)$$

$$\sigma_r^D = F_2^D - \frac{y^2}{Y_+} F_L^D$$

$$Y_+ = 1 + (1-y)^2$$

- the measurement will be performed in the same way as in case of F_L
- no direct measurement of F_L^D has been performed so far

$$F_L^D = \frac{Y_+^{460} Y_+^{920}}{y_{460}^2 Y_+^{920} - y_{920}^2 Y_+^{460}} (\sigma_r^{D920} - \sigma_r^{D460})$$

- x_{IP}, β, Q^2 fixed
- y vary $Q^2 = s x_{IP} \beta y$

Low Energy Runs at HERA

- last few weeks of HERA data taking were devoted to low energy runs that are necessary to measure the proton structure functions F_L and F_L^D
- at least 10 pb^{-1} of low energy data needed in order to perform a good F_L measurement
- $s = 4E_e E_p$
- 10 pb^{-1} of data with $E_p = 460 \text{ GeV}$ taken
- the low energy running was smooth, therefore the middle energy running had been proposed also
- 7 pb^{-1} of data with $E_p = 575 \text{ GeV}$ taken

Different Models

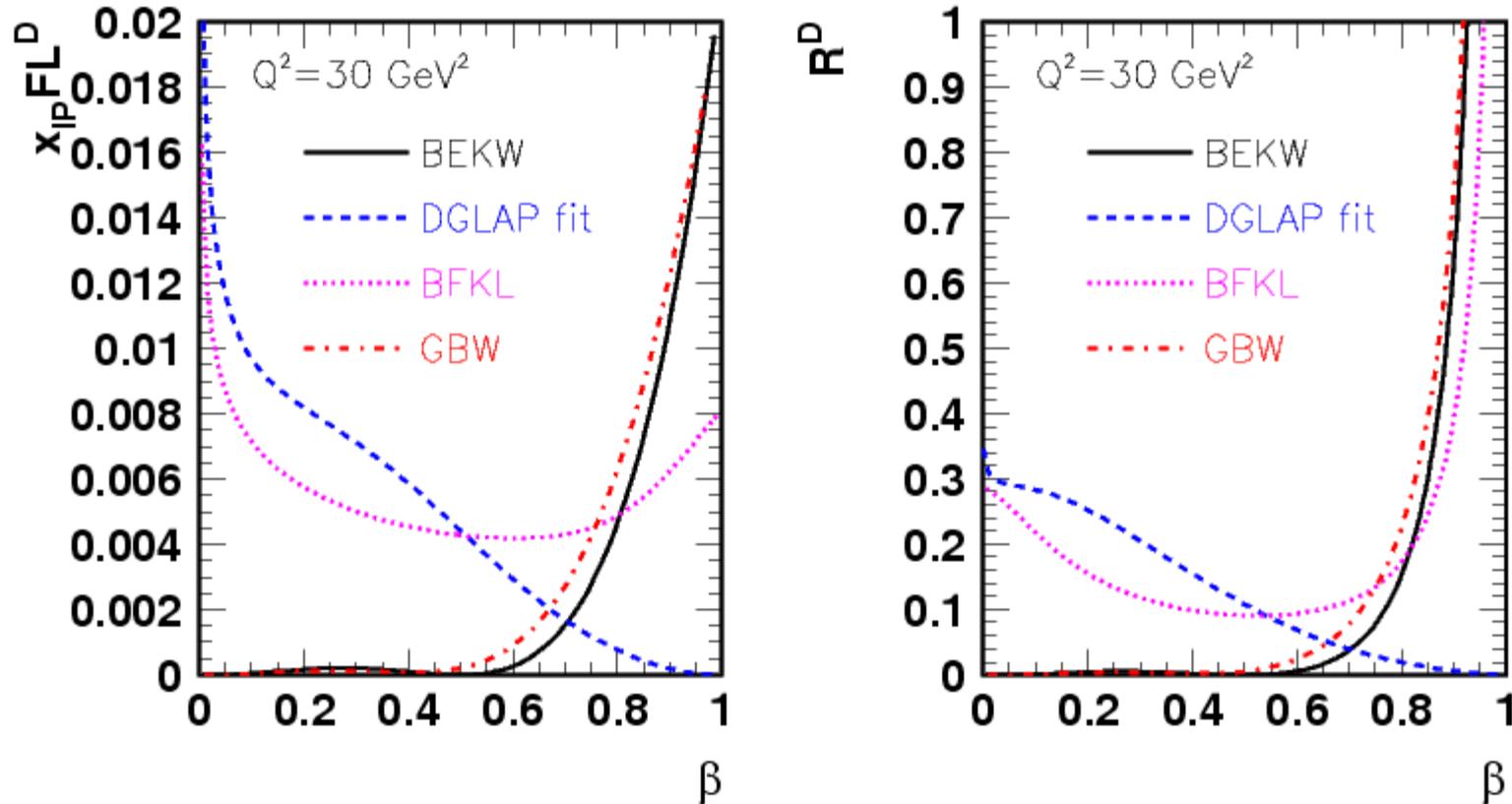


FIG. 9. Predictions for $x_{\mathbb{P}} F_L^D$ (left) and $R^D = \frac{F_L^D}{F_2^D - F_L^D}$ as a function of β at $Q^2 = 30 \text{ GeV}^2$ and $x_{\mathbb{P}} = 10^{-3}$. We present predictions for BEKW as a full line, DGLAP fit as a dashed line, BFKL as a dotted line and GBW as a dashed-dotted line.

Simulation

