

Searching for Hidden Dimensions in Space at DØ

Greg Landsberg

BROWN University

(For the  Collaboration)

Fermilab W&C Seminar

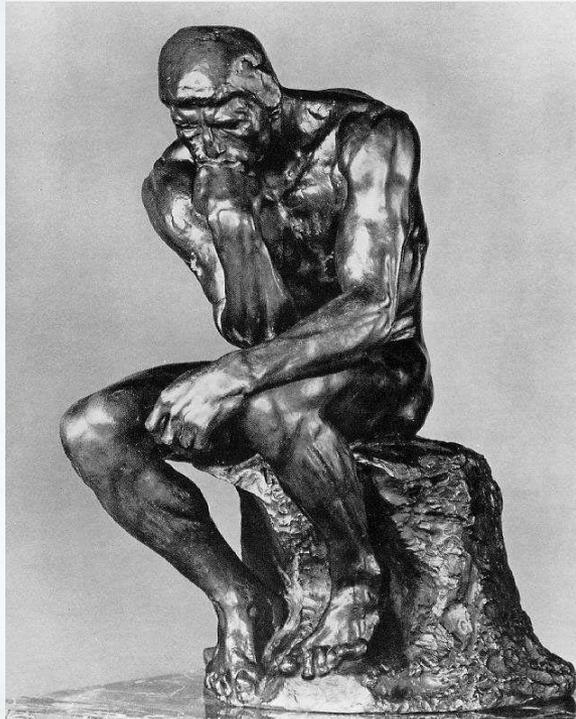
October 15, 2004





Out-Sketch

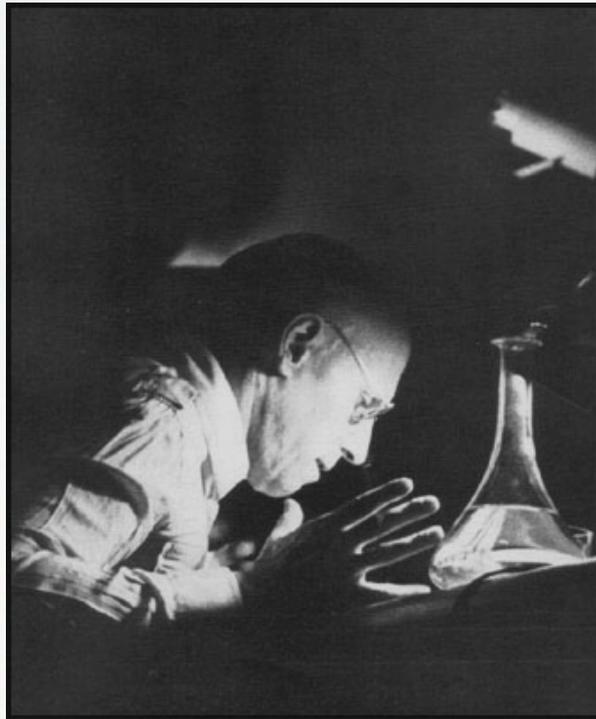
Theory/Phenomenology



Brief overview:

- ❖ Intro to Models
- ❖ Phenomenology
- ❖ Collider signatures

Tabletop Experiments



A few highlights:

- ❖ Gravity at short range
- ❖ Cosmology constraints

Accelerator Searches



Ultimate probes:

- ❖ New DØ Results on Searches for Extra Dimensions:

 - ❖ ADD Model
 - ❖ TeV⁻¹ Scenario
 - ❖ RS Model

- ❖ Outlook



N.B. Large Hierarchies Tend to Collapse...





Although Keep in Mind...

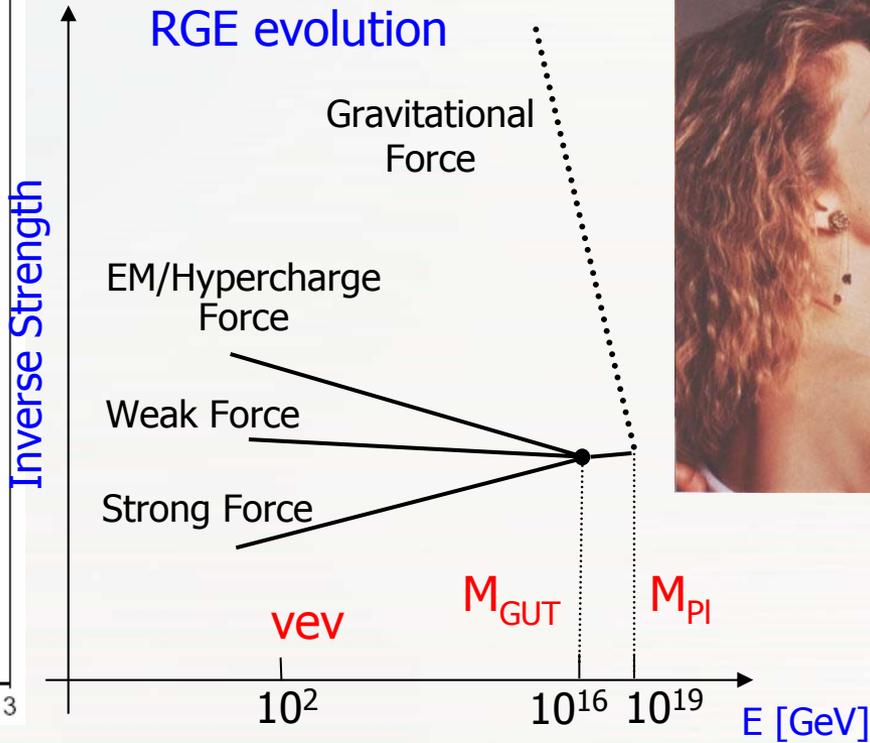
- **Fine tuning** (required to keep a large hierarchy stable) **exists in Nature**:
 - Solar eclipse: **angular size of the sun is the same as the angular size of the moon** within **2.5%** (pure coincidence!)
 - Numerology: **$987654321/123456789 = 8.000000073$**
(a bottle of wine is offered for the correct answer, why?)
 - Politics: **Florida recount, $2,913,321/2,913,144 = 1.000061$**
(a slice of cheese is offered for the correct answer, why?)
- [Answers should be given to me by the end of this talk]



Hierarchy of the Standard Model

Beauty ... and the Beast

	Measurement	Fit	$10^{\sigma_{meas}}$	$-10^{\sigma_{fit}}$
$\Delta\alpha_{had}^{(5)}(m_Z)$	0.02761 ± 0.00036	0.02768	0.1	0.1
m_Z [GeV]	91.1875 ± 0.0021	91.1873	0.1	0.1
Γ_Z [GeV]	2.4952 ± 0.0023	2.4965	0.1	0.1
σ_{had}^0 [nb]	41.540 ± 0.037	41.481	1.5	1.5
R_l	20.767 ± 0.025	20.739	0.5	0.5
$A_{tb}^{0,l}$	0.01714 ± 0.00095	0.01642	0.1	0.1
$A_l(P_Z)$	0.1465 ± 0.0032	0.1480	0.1	0.1
R_b	0.21638 ± 0.00066	0.21566	0.1	0.1
R_c	0.1720 ± 0.0030	0.1723	0.1	0.1
$A_{tb}^{0,b}$	0.0997 ± 0.0016	0.1037	2.5	2.5
$A_{tb}^{0,c}$	0.0706 ± 0.0035	0.0742	0.5	0.5
A_b	0.925 ± 0.020	0.935	0.5	0.5
A_c	0.670 ± 0.026	0.668	0.1	0.1
$A_l(\text{SLD})$	0.1513 ± 0.0021	0.1480	1.5	1.5
$\sin^2\theta_{eff}^{lept}(Q_{tb})$	0.2324 ± 0.0012	0.2314	0.1	0.1
m_W [GeV]	80.425 ± 0.034	80.398	0.1	0.1
Γ_W [GeV]	2.133 ± 0.069	2.094	0.5	0.5
m_t [GeV]	178.0 ± 4.3	178.1	0.1	0.1



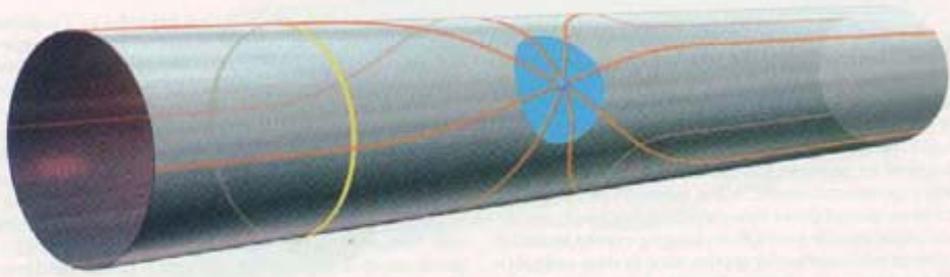
Extra dimensions might get rid of the beast while preserving the beauty!

The ADD Model

- Arkani-Hamed, Dimopoulos, and Dvali, 1998 (ADD)
- SM fields are localized on the (3+1)-brane; gravity is the only gauge force that “feels” the bulk space
- What about Newton’s law?

$$V(r) = \frac{1}{M_{Pl}^2} \frac{m_1 m_2}{r} \rightarrow \frac{1}{(M_{Pl}^{[3+n]})^{n+2}} \frac{m_1 m_2}{r^{n+1}}$$

- Ruled out for infinite extra dimensions, but does not apply for sufficiently small compactified extra dimensions:



$$V(r) \propto \frac{1}{(M_{Pl}^{[3+n]})^{n+2}} \frac{m_1 m_2}{R^n r} \text{ for } r \gg R$$

- ✚ Gravity is fundamentally strong force, but we do not feel that as it is diluted by the volume of the bulk
- ✚ $G'_N = 1/(M_{Pl}^{[3+n]})^2 \equiv 1/M_D^2$; $M_D \sim 1 \text{ TeV}$

$$M_D^{n+2} \propto M_{Pl}^2 / R^n$$
- ✚ More precisely, from Gauss’s law:

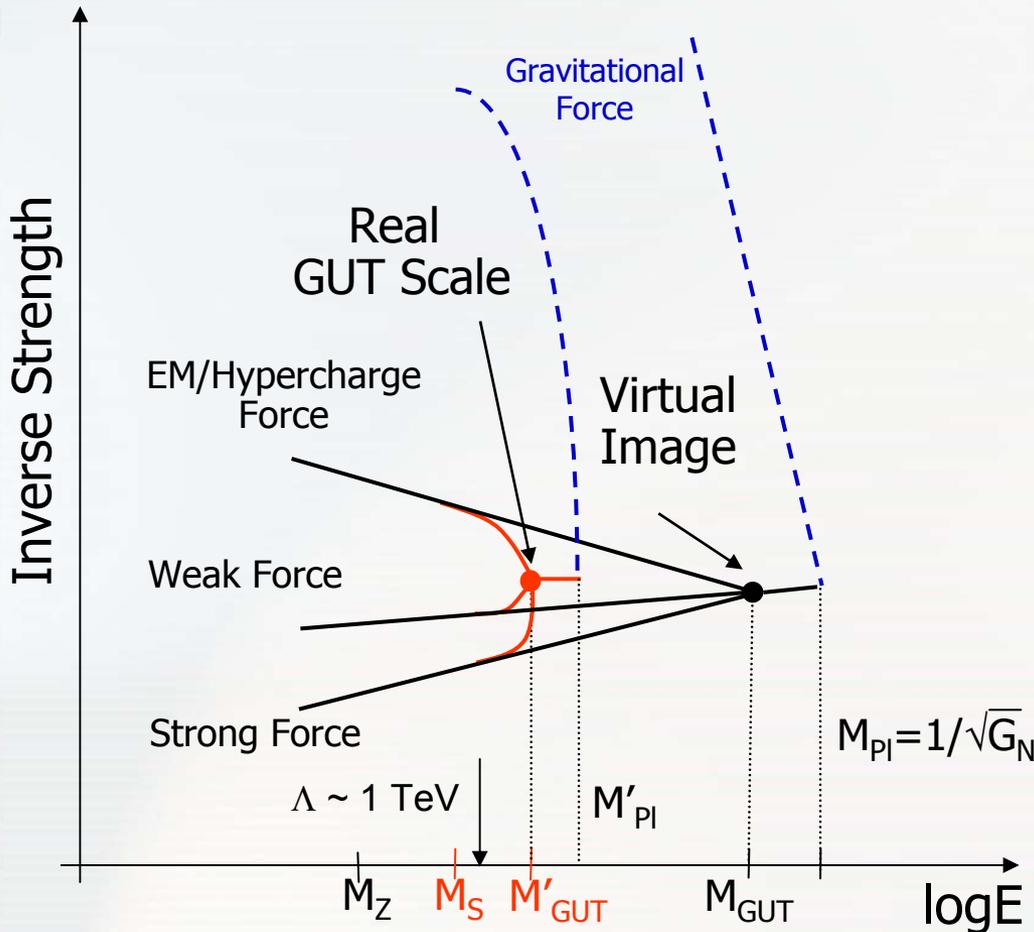
$$R = \frac{1}{2\sqrt{\pi} M_D} \left(\frac{M_{Pl}}{M_D} \right)^{2/n} \propto \begin{cases} 8 \times 10^{12} m, & n = 1 \\ 0.7 \text{ mm}, & n = 2 \\ 3 \text{ nm}, & n = 3 \\ 6 \times 10^{-12} m, & n = 4 \end{cases}$$

- ✚ Amazing as it is, but no one has tested Newton’s law to distances less than $\sim 1 \text{ mm}$ (as of 1998)
- ✚ Thus, the fundamental Planck scale could be as low as 1 TeV for $n > 1$



Longitudinal ED

• Simultaneously, another idea has appeared:



– Explore modification of the RGE in $(4+n)$ -dimensions to **achieve low-energy unification of three gauge forces** [Dienes, Dudas, and Gherghetta, PL B436, 55 (1998)]

– To achieve that, **allow gauge bosons** (g , γ , W , and Z) **to propagate in an extra dimension**, which is “longitudinal” to the SM brane and compactified on a “natural” EW scale: **$R \sim 1 \text{ TeV}^{-1}$**



Randall-Sundrum Scenario

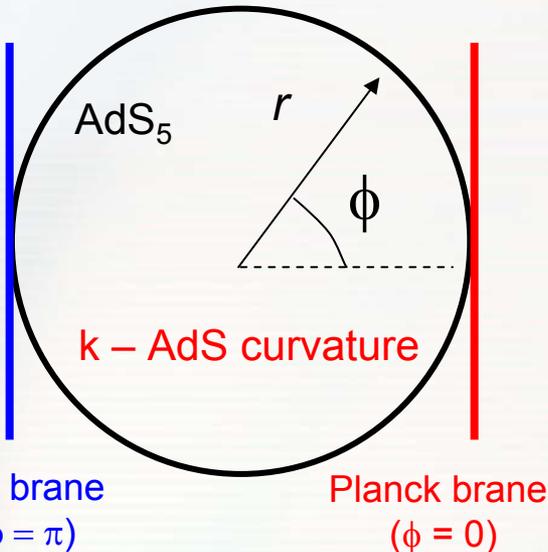
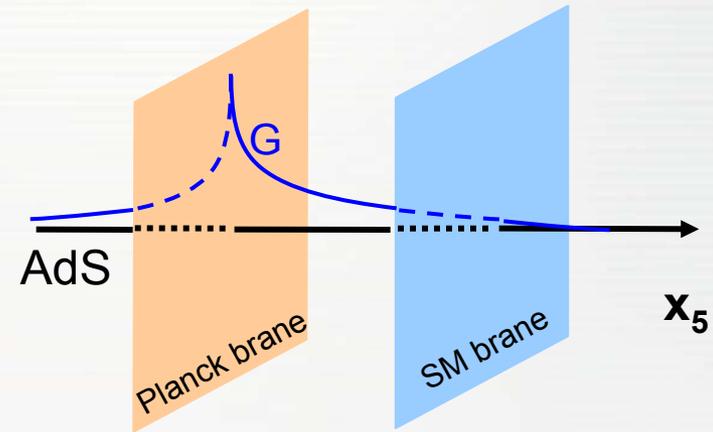
• Randall-Sundrum (RS) scenario
 [PRL **83**, 3370 (1999); PRL **83**, 4690 (1999)]

❖ + brane – no low energy effects

❖ +– branes – TeV Kaluza-Klein modes of graviton

❖ Low energy effects on SM brane are given by Λ_π ;

❖ for $kr_c \sim 10$, $\Lambda_\pi \sim 1$ TeV and the hierarchy problem is solved naturally



$$ds^2 = e^{-2kr|\phi|} \eta_{\mu\nu} dx^\mu dx^\nu - r^2 d\phi^2$$

$$\Lambda_\pi = \bar{M}_{Pl} e^{-kr\pi}$$

Reduced Planck mass:

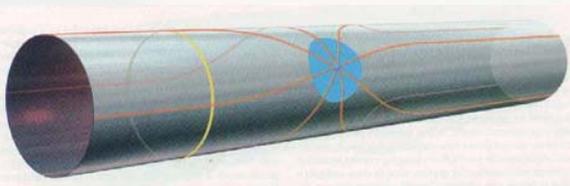
$$\bar{M}_{Pl} \equiv M_{Pl} / \sqrt{8\pi}$$



Differences Between the Models

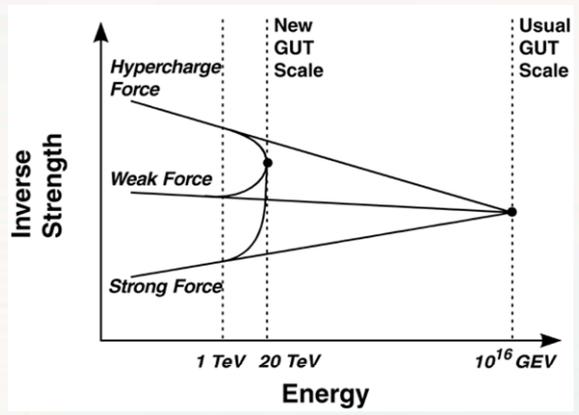
ADD Model:

- Pro: “Eliminates” the hierarchy problem by stating that physics ends at a TeV scale
- Only gravity lives in the “bulk” space
- Size of ED’s (n=2-7) between $\sim 100 \mu\text{m}$ and $\sim 1 \text{ fm}$
- Black holes at the LHC
- Con: Doesn’t explain how to make ED large



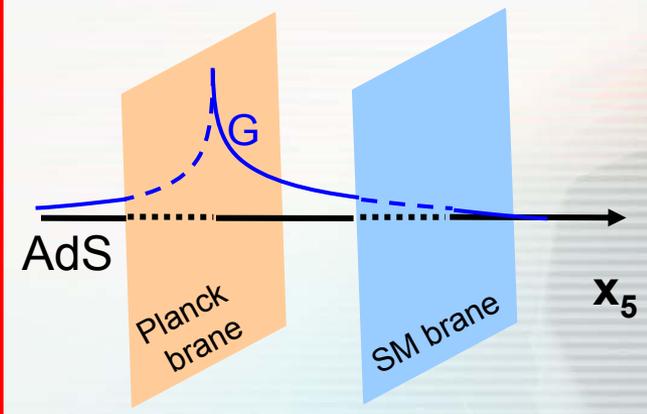
TeV⁻¹ Scenario:

- Pro: Lowers GUT scale by changing running of the couplings
- Only gauge bosons (g/γ/W/Z) “live” in ED’s
- Size of ED’s $\sim 1 \text{ TeV}^{-1}$ or $\sim 10^{-19} \text{ m}$
- Con: Gravity is not in the picture



RS Model:

- Pro: A rigorous solution to the hierarchy problem via localization of gravity
- Gravitons (and possibly other particles) propagate in a single ED, w/ special metric
- Con: Size of ED as small as $\sim 1/M_{\text{Pl}}$ or $\sim 10^{-35} \text{ m}$

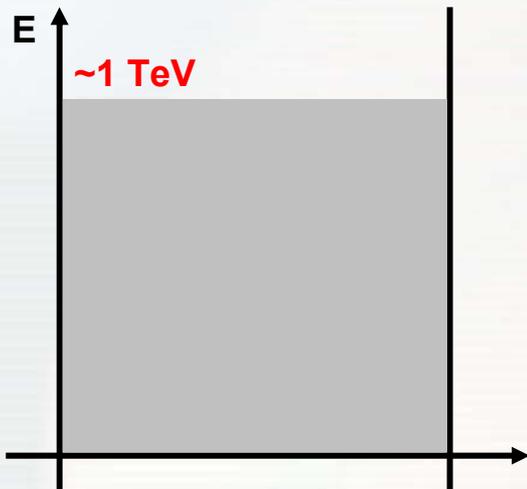




Kaluza-Klein Spectrum

ADD Model:

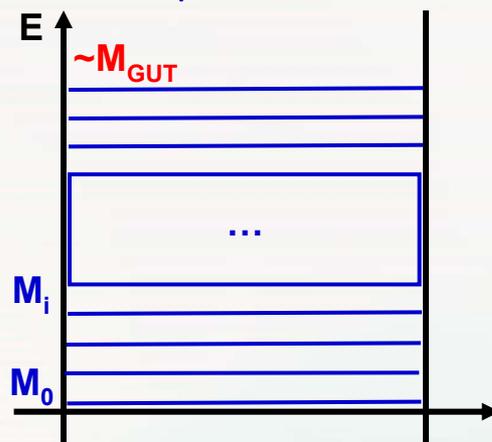
- Winding modes with energy spacing $\sim 1/r$, i.e. 1 meV – 100 MeV
- Can't resolve these modes – they appear as continuous spectrum
- Coupling: G_N per mode; compensated by large number of modes



TeV⁻¹ Scenario:

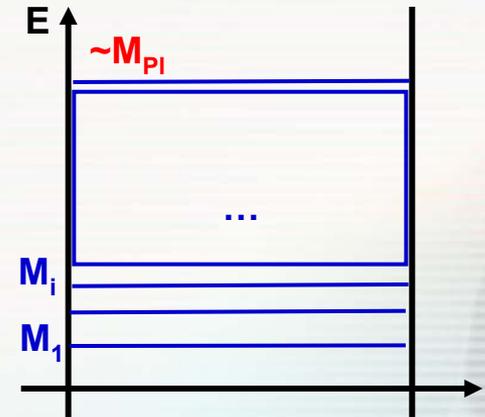
- Winding modes with nearly equal energy spacing $\sim 1/r$, i.e. $\sim \text{TeV}$
- Can excite individual modes at colliders or look for indirect effects
- Coupling: $\sim g_w$ per mode

$$M_i = \sqrt{M_0^2 + i^2/r^2}$$



RS Model:

- “Particle in a box” with special AdS metric
- Energy eigenvalues are given by zeroes of Bessel function J_1
- Light modes might be accessible at colliders
- Coupling: G_N for zero mode; $1/\Lambda_\pi^2$ for the others

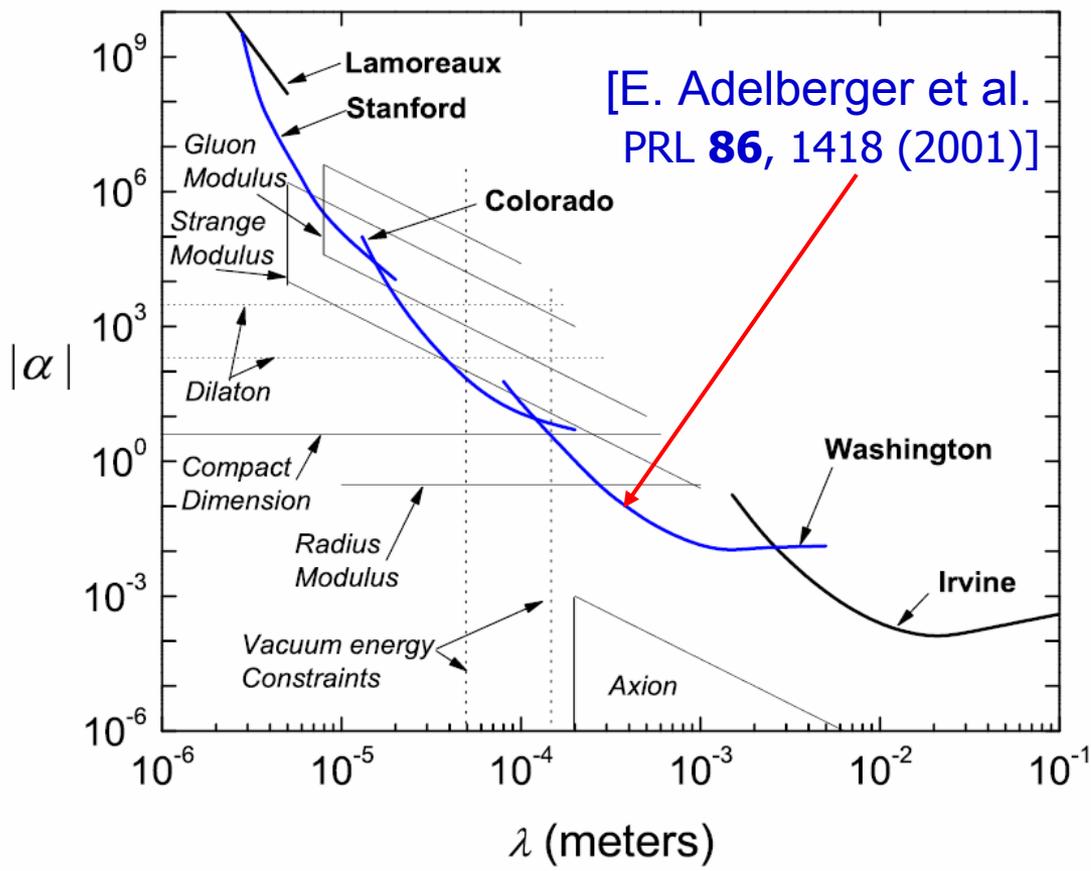


$$M_0 = 0; M_i = M_1 x_i/x_1 \approx M_1, 1.83M_1, 2.66M_1, 3.48M_1, 4.30M_1, \dots$$



Constraints from Gravity Experiments

[J. Long and J. Price, hep-ph/0303057]



- **Sub-millimeter gravity measurements** could probe *only* $n=2$ case *only* within the ADD model

- The **best sensitivity** so far have been achieved in the **U of Washington torsion balance experiment** – a high-tech “remake” of the 1798 Cavendish experiment

❖ $R \lesssim 0.16 \text{ mm} (M_D \gtrsim 1.7 \text{ TeV})$

- **Sensitivity vanishes quickly with the distance** – can’t push limits further down significantly

- Started restricting **ADD with 2 extra dimensions**; can’t probe any higher number

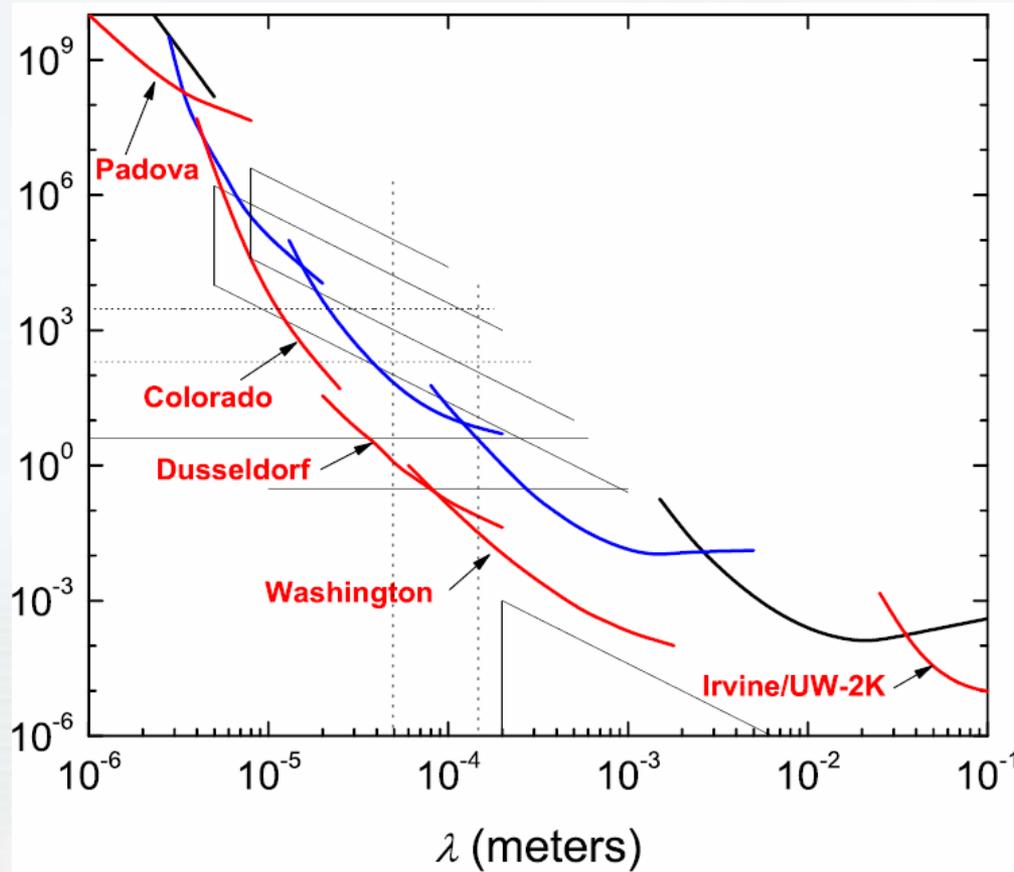
- Ultimately **push the sensitivity by a factor of two** in terms of the distance

- **No sensitivity** to the TeV^{-1} and RS models



Constraints from Gravity Experiments

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Astrophysical and Cosmological Constraints

- Used to constrain the **ADD scenario**
- **Supernova cooling due to graviton emission – an alternative cooling mechanism** that would **decrease** the dominant cooling via **neutrino emission**
 - ❖ Tightest **limits** on any additional cooling sources come **from the measurement of the SN1987A neutrino flux** by the Kamiokande and IMB
 - ❖ Application to the ADD scenario [Cullen and Perelstein, PRL **83**, 268 (1999); Hanhart, Phillips, Reddy, and Savage, NP **B595**, 335 (2001)]:
 - ❖ $M_D > 25\text{-}30 \text{ TeV}$ ($n=2$)
 - ❖ $M_D > 2\text{-}4 \text{ TeV}$ ($n=3$)
- **Distortion of the cosmic diffuse gamma radiation (CDG) spectrum** due to the G_{KK}
→ $\gamma\gamma$ decays [Hall and Smith, PRD **60**, 085008 (1999)]:
 - ❖ $M_D > 100 \text{ TeV}$ ($n=2$)
 - ❖ $M_D > 5 \text{ TeV}$ ($n=3$)
- **Overclosure of the universe, matter dominance in the early universe** [Fairbairn, PL **B508**, 335 (2001); Fairbairn and Griffiths, JHEP **0202**, 024 (2002)]
 - $M_D > 86 \text{ TeV}$ ($n=2$)
 - $M_D > 7.4 \text{ TeV}$ ($n=3$)
- **Neutron star γ -emission from radiative decays of the gravitons** trapped during the supernova collapse [Hannestad and Raffelt, PRL **88**, 071301 (2002)]:
 - $M_D > 1700 \text{ TeV}$ ($n=2$)
 - $M_D > 60 \text{ TeV}$ ($n=3$)
- **Caveat:** there are **many** known (and unknown!) **uncertainties**, so the cosmological bounds are reliable **only as an order of magnitude estimate**
- Still, **$n=2$** is largely **disfavored**
- **Virtually no constraints on TeV^{-1} and RS models**

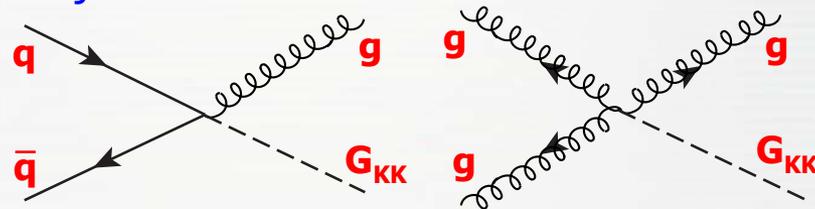


Collider Signatures for Large ED

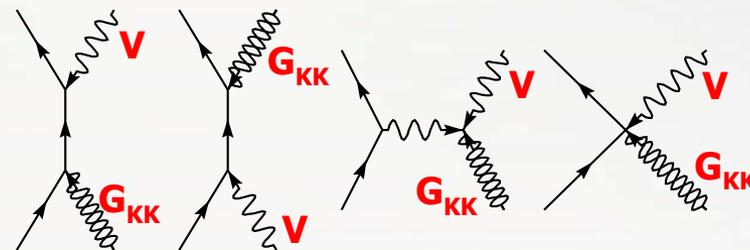
- Kaluza-Klein gravitons couple to the energy-momentum tensor, and therefore contribute to most of the SM processes
- For Feynman rules for G_{KK} see:
 - [Han, Lykken, Zhang, PRD **59**, 105006 (1999)]
 - [Giudice, Rattazzi, Wells, NP **B544**, 3 (1999)]
- Since graviton can propagate in the bulk, energy and momentum are not conserved in the G_{KK} emission from the point of view of our 3+1 space-time
- Depending on whether the G_{KK} leaves our world or remains virtual, the collider signatures include single photons/Z/jets with missing E_T or fermion/vector boson pair production
- Graviton emission: direct sensitivity to the fundamental Planck scale M_D
- Virtual effects: sensitive to the ultraviolet cutoff M_S , expected to be $\sim M_D$ (and likely $< M_D$)
- The two processes are complementary

Real Graviton Emission

Monojets at hadron colliders

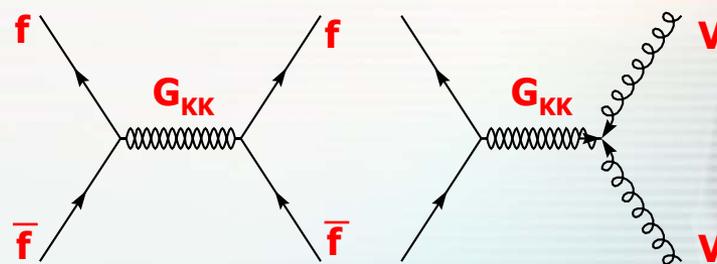


Single VB at hadron or e^+e^- colliders



Virtual Graviton Effects

Fermion or VB pairs at hadron or e^+e^- colliders





L'Épilogue

Experiment	$e^+e^- \rightarrow \gamma G$					$e^+e^- \rightarrow ZG$					Color coding
	n=2	n=3	n=4	n=5	n=6	n=2	n=3	n=4	n=5	n=6	
ALEPH	1.28	0.97	0.78	0.66	0.57	0.35	0.22	0.17	0.14	0.12	≤ 184 GeV
DELPHI	1.38	1.02	0.84	0.68	0.58	 	 	 	 	 	≤ 189 GeV
L3	1.02	0.81	0.67	0.58	0.51	0.60	0.38	0.29	0.24	0.21	> 200 GeV
OPAL	1.09	0.86	0.71	0.61	0.53	 	 	 	 	 	$\lambda=-1$ $\lambda=+1$

Virtual Graviton Exchange [M_s(Hewett)]

Experiment	e^+e^-	$\mu^+\mu^-$	$\tau^+\tau^-$	qq	ff	$\gamma\gamma$	WW	ZZ	Combined
ALEPH	1.04 0.81	0.65 0.67	0.60 0.62	0.53/0.57 0.46/0.46 (bb)	1.05 0.84	0.81 0.82	 	 	0.75/1.00 (<189)
DELPHI	 	0.59 0.73	0.56 0.65	 	0.60 0.76	0.83 0.91	 	 	0.60/0.76 (ff) (<202)
L3	0.98 1.06	0.56 0.69	0.58 0.54	0.49 0.49	0.84 1.00	0.99 0.84	0.68 0.79	 	1.1/1.0 (<202)
OPAL	1.15 1.00	0.62 0.66	 	 	0.62 0.66	0.89 0.83	 	0.63 0.74	1.17/1.03 (<209)

LEP Combined: 1.2/1.1

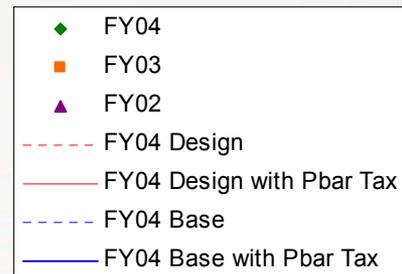
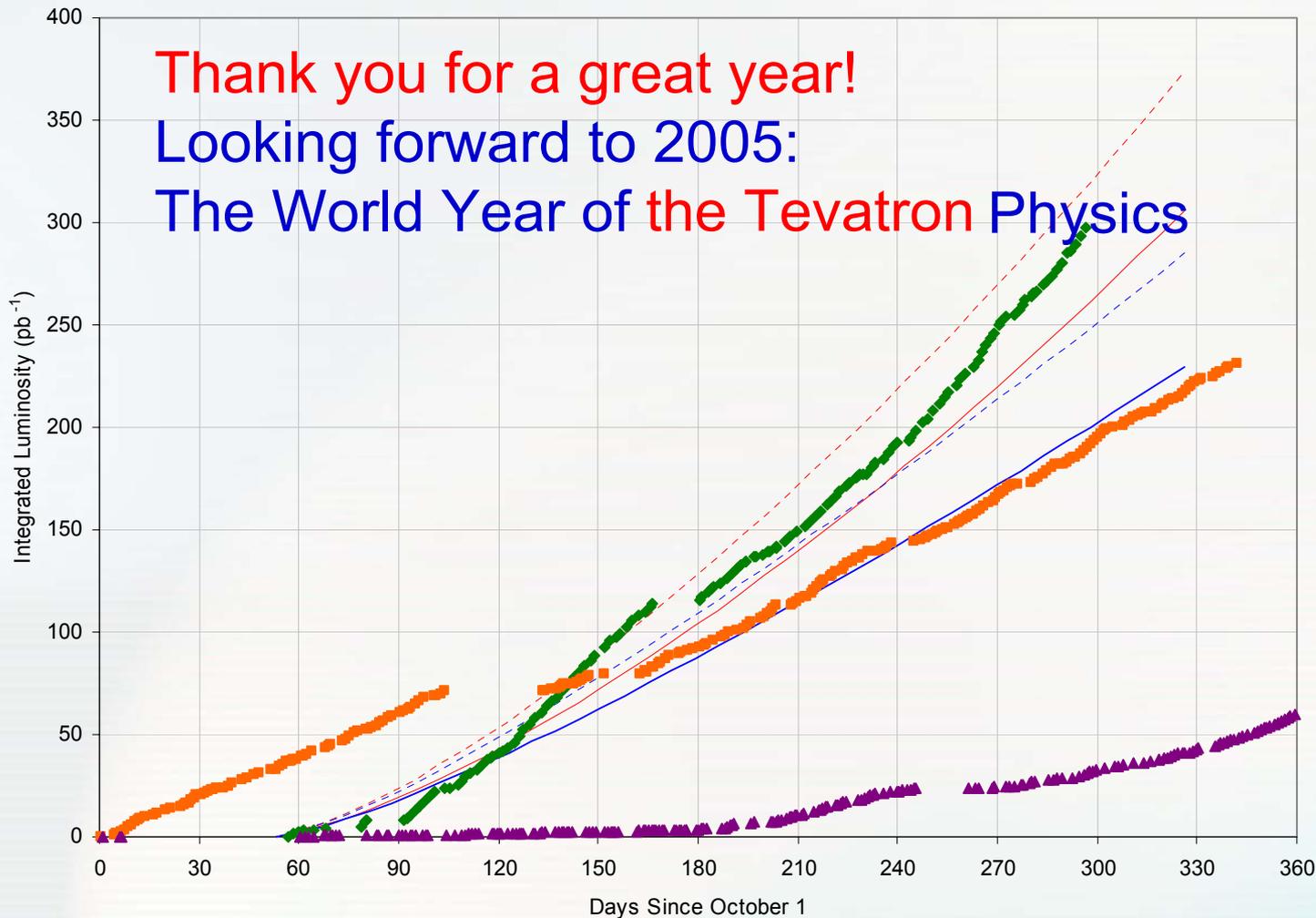


The Tevatron Performance

The FY2004 Design goal (306 pb⁻¹ delivered) has been achieved

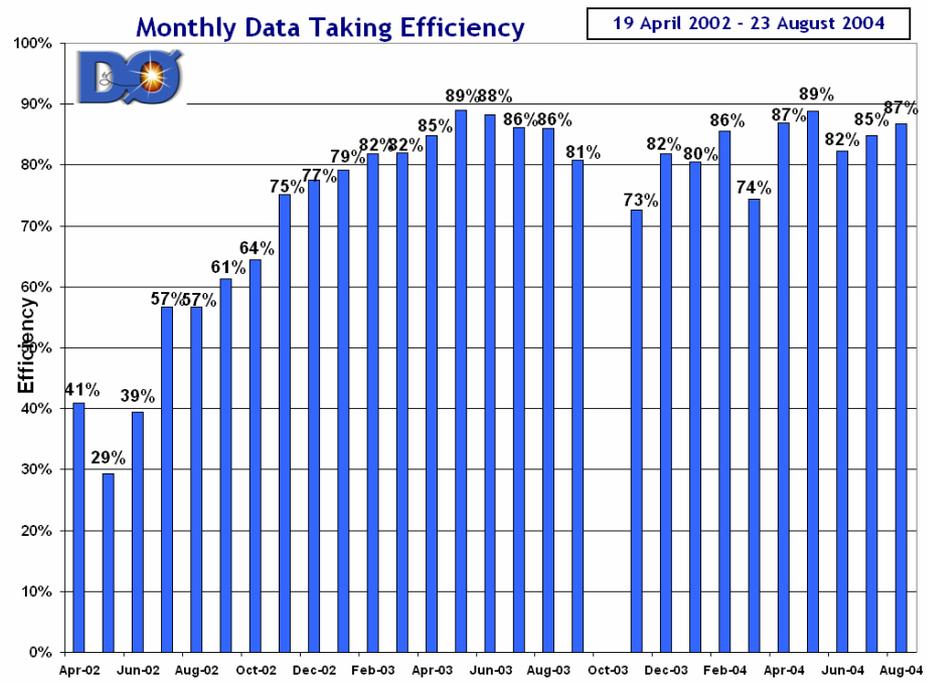
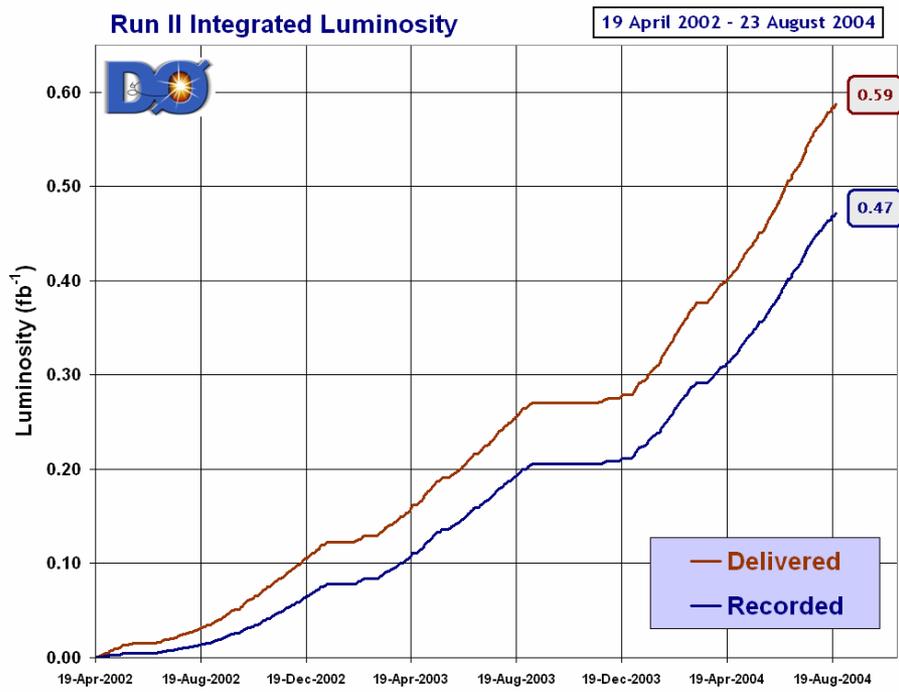
$L = 10^{32} \text{ cm}^{-2}\text{s}^{-1}$
Record!

Thank you for a great year!
Looking forward to 2005:
The World Year of the Tevatron Physics



Nearly 0.5 fb^{-1} has been logged to tape

1 billion events recorded as of August 2!



- ❖ The experiment is $\approx 85\%$ efficient over the last year; 80% average efficiency in Run II
- ❖ Most of the analyses shown in this talk are based on $\sim 200 \text{ pb}^{-1}$, which corresponds to the 2002-2003 data



Monojets: Tainted History

EXPERIMENTAL OBSERVATION OF EVENTS WITH LARGE MISSING TRANSVERSE ENERGY
ACCOMPANIED BY A JET OR A PHOTON(S) IN $p\bar{p}$ COLLISIONS

AT $\sqrt{s} = 540$ GeV

[PL, 139B, 115 (1984)]

UA1 Collaboration, CERN, Geneva, Switzerland

VOLUME 54, NUMBER 6

PHYSICAL REVIEW LETTERS

11 FEBRUARY 1985

Monojets from Z Decay without Extra Neutrinos or Higgs Particles

Stephen F. King

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138

(Received 26 November 1984)

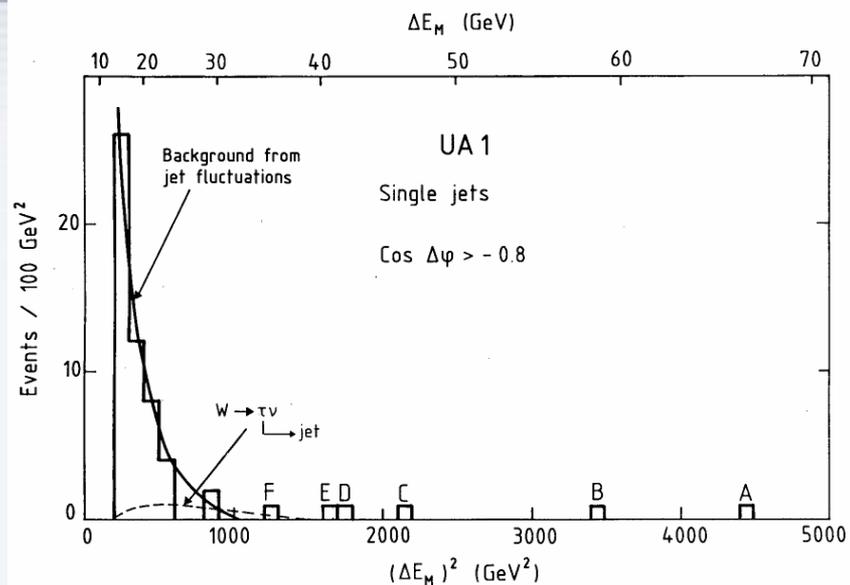
closely spaced photons). We cannot find an explanation for such events in terms of backgrounds or within the expectations of the Standard Model.

The recent discovery of monojets by Arnison *et al.*¹ at the CERN $p\bar{p}$ collider has caused ripples of excitement throughout the particle physics world, since they cannot be explained by the minimal standard model.²

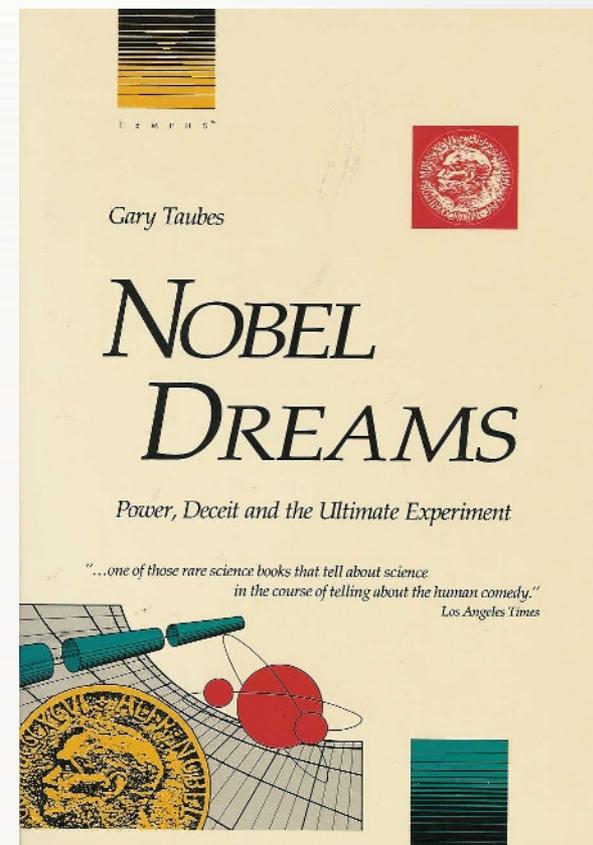




Monojets: Tainted History



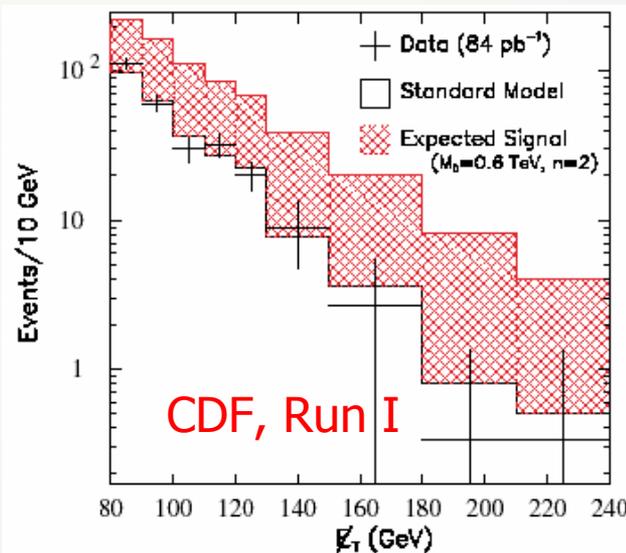
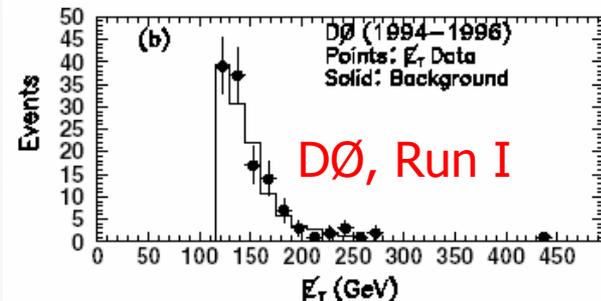
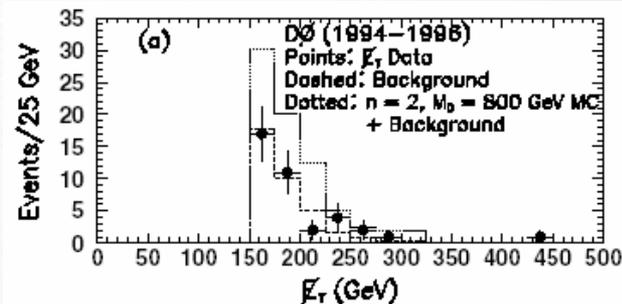
- These **monojets** turned out to be due to **unaccounted background**
- The **signature was deemed doomed** and nearly forgotten
- It **took many years for the successful monojet analyses** at a hadron collider to be completed (at CDF and DØ)





Search for Monojets at the Tevatron

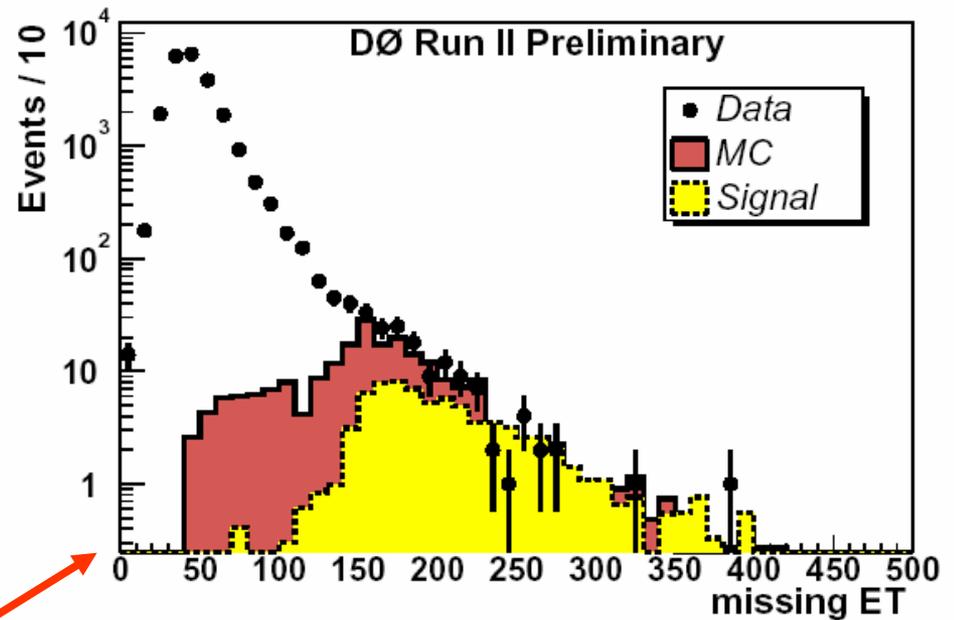
- Challenge: large instrumental background from mismeasured ME_T , cosmics, and $W(\tau\nu)$
- Irreducible physics background from $Z(\nu\nu) + \text{jet}(s)$: forces to use high jet p_T , ME_T cuts
- Pioneered in Run I by DØ
[PRL 90, 251802 (2003)]:
 $M_D > 0.63 - 0.89$ TeV ($n=6-2$)
- Recently superseded by CDF
[PRL 92, 121802 (2004)]:
 $M_D > 0.71 - 1.00$ TeV ($n=6-2$)
- CDF has pioneered similar search in $\gamma + ME_T$, albeit less sensitive
[PRL 89, 281801 (2002)]





Monojet Event Selection

- New Run II DØ analysis based on $\sim 85 \text{ pb}^{-1}$ of data, collected with special trigger, requiring $M_{H_T} = -\sum \vec{p}_T(\text{jets}) > 20 \text{ GeV}$
- Offline, jets require to have associated tracks and confirm primary vertex
- Jet quality cuts reduce instrumental background from calorimeter noise
- Leading jet $p_T > 150 \text{ GeV}$;
- $M_{E_T} > 150 \text{ GeV}$
- No extra jets with $p_T > 50 \text{ GeV}$ are allowed; no isolated leptons
- M_{E_T} vector is required to be separated by at least 30° in azimuth from any jet in the event

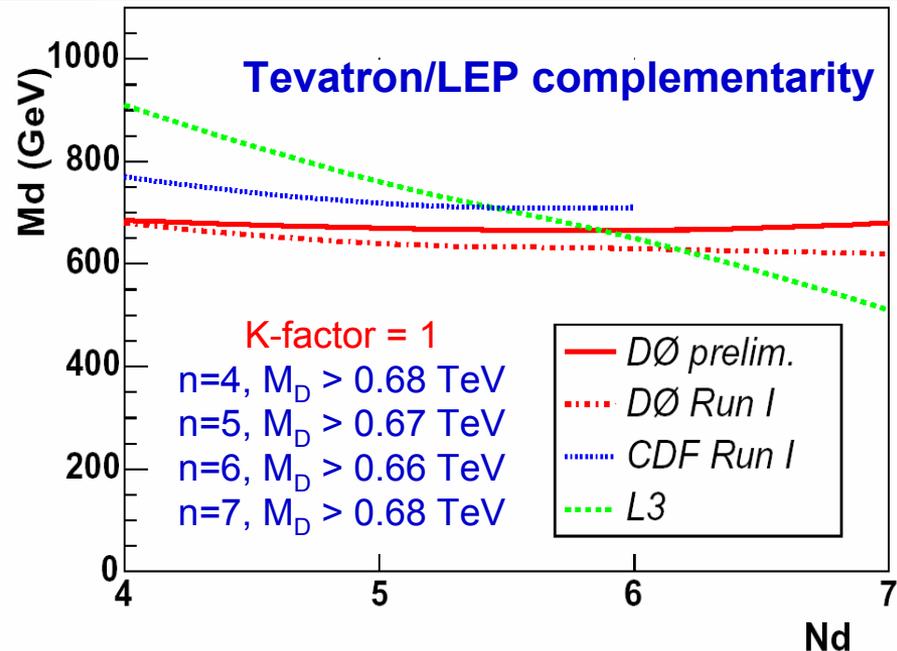
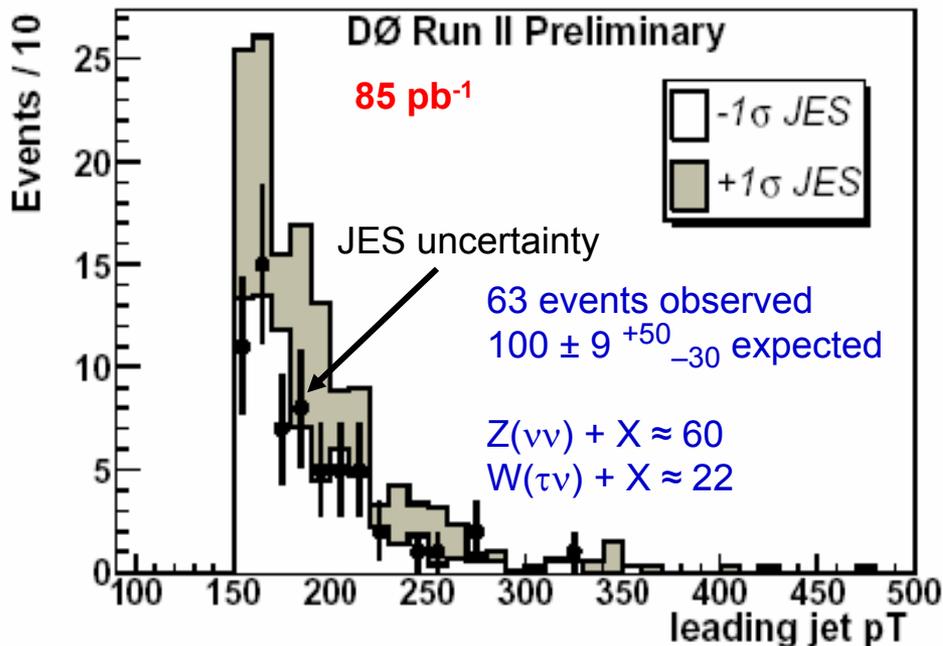


- Overall signal ($M_D = 0.7 \text{ TeV}$, $n=6$) efficiency: $\sim 5\%$; mainly due to high p_T cut
- 63 events (out of the initial sample of 360K) pass final selections
- Signal simulation: PYTHIA augmented with Lykken-Matchev's code



Monojet Results

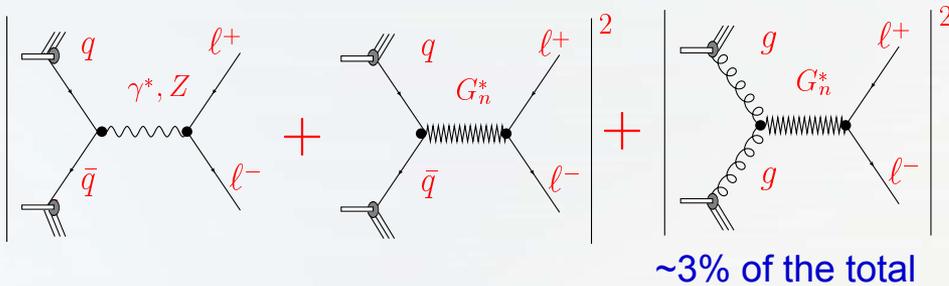
- Major systematics from **jet energy scale** – to be reduced soon (JES errors have been reduced by a factor of ~ 2 since these results became public)
- **Sensitivity** already **exceeds that for DØ in Run I**, but **still below CDF's Run I result**
- Impressive sensitivity already **achieved with less data** due to superior detector and higher energy
- **An update** with much higher statistics and better detector understanding **soon**





Virtual Graviton Effects

• In the case of **pair production** via virtual graviton, **gravity effects interfere with the SM** (e.g., l^+l^- at hadron colliders):



• Therefore, **production cross section has three terms**: SM, interference, and direct gravity effects:

$$\frac{d^2\sigma}{d\cos\theta^*dM} = \frac{d^2\sigma_{SM}}{d\cos\theta^*dM} + \eta_G f_1(\cos\theta^*, M) + \eta_G^2 f_2(\cos\theta^*, M)$$

$$\eta_G = a(n)/M_S^4, \quad a(n) \sim 1$$

• The **sum in KK states is divergent** in the effective theory, so in order to calculate the cross sections, **an explicit cut-off is required**

• An expected value of the **cut-off $M_S \approx M_D$** , as this is the scale at which the effective theory breaks down, and the string theory needs to be used to calculate production

• There are **three major conventions** on how to write the **effective Lagrangian**:

- Hewett [PRL **82**, 4765 (1999)]
- Giudice, Rattazzi, Wells [NP **B544**, 3 (1999); revised version, hep-ph/9811291]
- Han, Lykken, Zhang [PRD **59**, 105006 (1999); revised version, hep-ph/9811350]

• Fortunately **all three conventions** turned out to be **equivalent** and only the **definitions of M_S are different**



Hewett, GRW, and HLZ Formalisms

- **Hewett**: neither sign of the interference nor the dependence on the number of extra dimensions is known; therefore the **interference term is $\sim \lambda/M_S^4$ (Hewett)**, where λ is of order 1; numerically uses $\lambda = \pm 1$
- **GRW**: sign of the interference is fixed, but the dependence on the number of extra dimensions is unknown; therefore the **interference term is $\sim 1/\Lambda_T^4$** (where Λ_T is their notation for M_S)
- **HLZ**: not only the sign of interference is fixed, but the n-dependence can be calculated in the effective theory; thus the **interference term $\sim a(n)/M_S^4$ (HLZ)**, where $a(n)$ reflects the dependence on the number of extra dimensions:

$$a(n) = \begin{cases} \log\left(\frac{M_S^2}{s}\right), & n = 2 \\ \frac{2}{n-2}, & n > 2 \end{cases}$$

- **Correspondence** between the three formalisms:

$$M_S(\mathbf{Hewett})|_{\lambda=\pm 1} \equiv \sqrt[4]{\frac{2}{\pi}} \Lambda_T(\mathbf{GRW})$$

$$\frac{\lambda}{M_S^4(\mathbf{Hewett})} = \frac{\pi}{2} \frac{\mathcal{F}}{M_S^4(\mathbf{HLZ})}$$

$$\frac{1}{\Lambda_T^4(\mathbf{GRW})} = \frac{\mathcal{F}}{M_S^4(\mathbf{HLZ})}$$

- **Rule of thumb**:

$$M_S(\mathbf{Hewett})|_{\lambda=\pm 1} \approx M_S(\mathbf{HLZ})|_{n=5}$$

$$\Lambda_T(\mathbf{GRW}) = M_S(\mathbf{HLZ})|_{n=4}$$



Search for Large ED in Dimuons

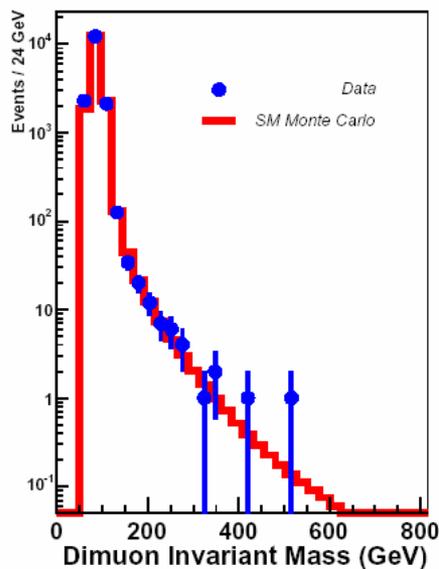
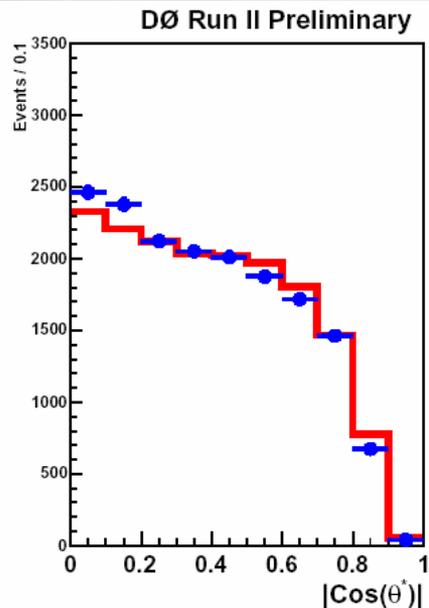
- Data set: $250 \pm 16 \text{ pb}^{-1}$, $\sim 17\text{K}$ dimuon candidate events:

Selection	Number of events passing cut
Starting sample	115,009
Bad run removal	108,574
Duplicate events removed	105,863
Track quality cuts	65,163
Dimuon invariant mass $> 50 \text{ GeV}$	40,744
Cosmic veto	25,811
Isolation requirements	17,193
After p_T fixed re-apply mass cut	16,796

- QCD and cosmic backgrounds made negligible by tight muon ID cuts (isolation and timing)
- Major remaining background: SM Drell-Yan production
- Use special technique to take care of poorly measured high- p_T muons
- Use 2D-analysis based on $M_{\mu\mu}$ and $|\cos\theta^*|$ yielding maximum sensitivity to large ED [Cheung & GL, PRD **62**, 076003 (2000)]
- Extract signal by fitting data with the sum of the SM, interference, and direct graviton templates, with η_G being a parameter of the fit



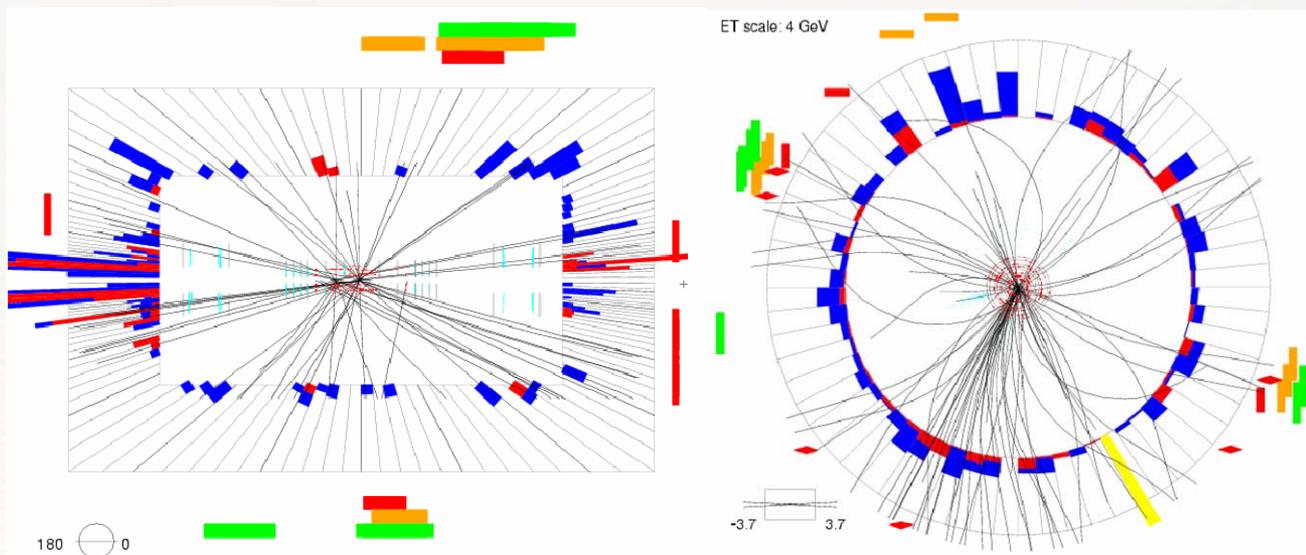
Comparison of Data and the SM



Min. $M_{\mu\mu}$	Bckgd	N in data	Poisson P
120 GeV	221	213	0.71
180 GeV	43.6	43	0.56
240 GeV	15.1	15	0.54
300 GeV	6.4	5	0.45
360 GeV	3.2	2	0.82
420 GeV	1.7	1	0.82
480 GeV	1.0	1	0.64
510 GeV	0.83	1	0.56
540 GeV	0.67	0	1.0

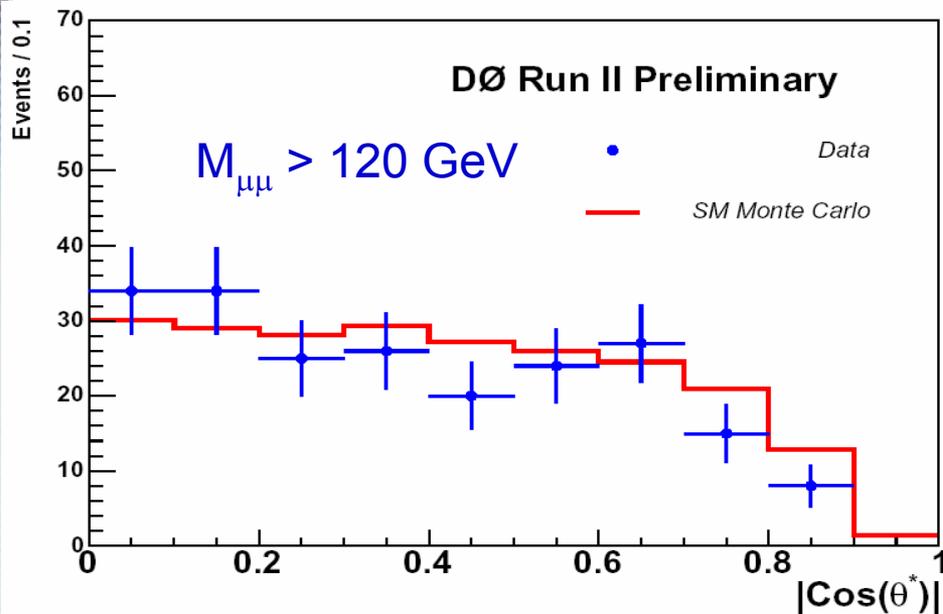
The highest-mass candidate event:

- ❖ $M_{\mu\mu} = 522$ GeV
- ❖ $|\text{cos}\theta^*| = 0.15$
- ❖ $N_{\text{jet}} (p_T > 15 \text{ GeV}) = 0$
- ❖ $M E_T^{\text{cal}} = 6$ GeV
- ❖ $p_T = 258$ GeV
- ❖ $\eta_{1,2} = 0.03, 0.4$
- ❖ $\phi_{1,2} = 5.9, 2.6$





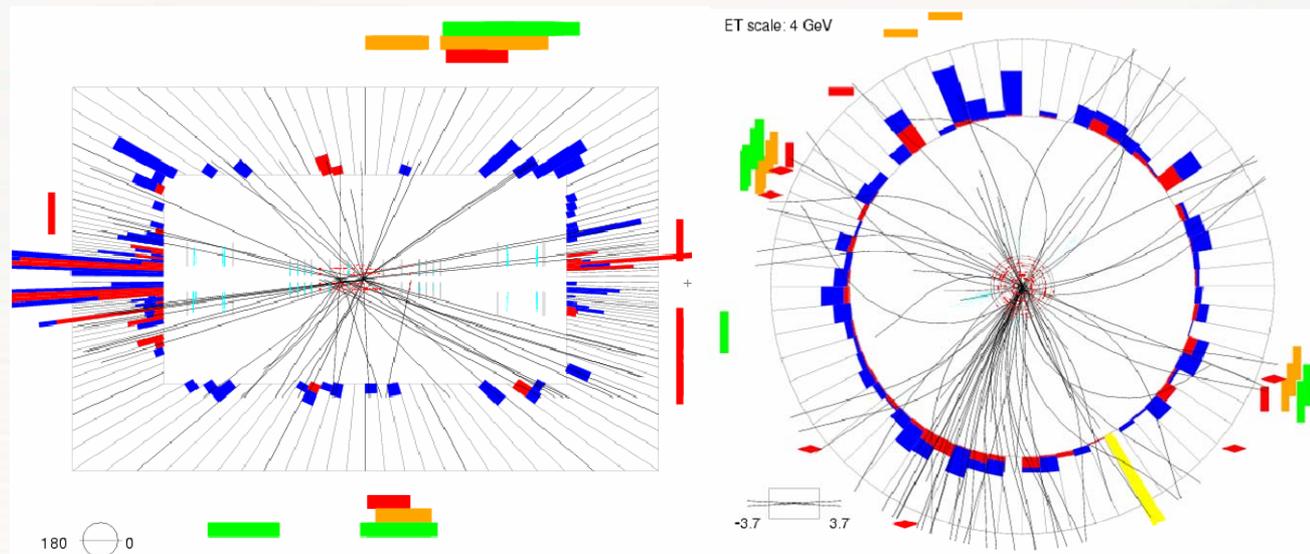
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The highest-mass candidate event:

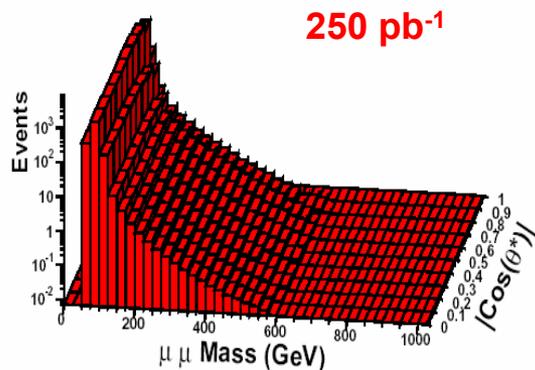
- ❖ $M_{\mu\mu} = 522 \text{ GeV}$
- ❖ $|\cos\theta^*| = 0.15$
- ❖ $N_{\text{jet}} (p_T > 15 \text{ GeV}) = 0$
- ❖ $M E_T^{\text{cal}} = 6 \text{ GeV}$
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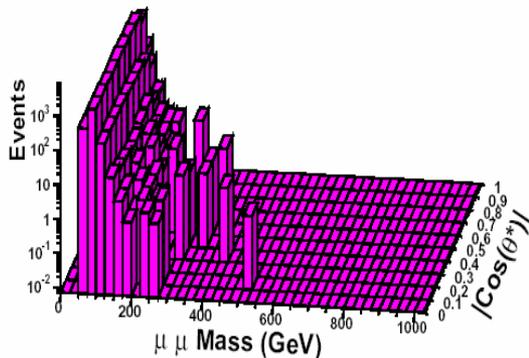
Limits on Large ED from Dimuons

Standard Model Monte Carlo



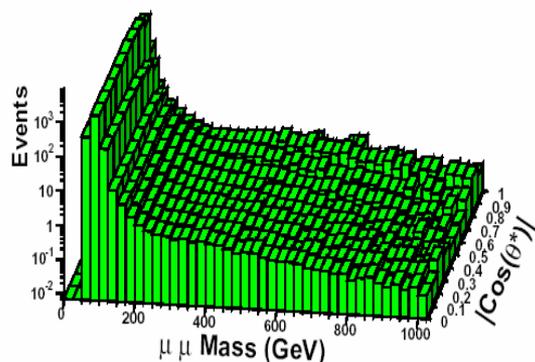
250 pb⁻¹

Data

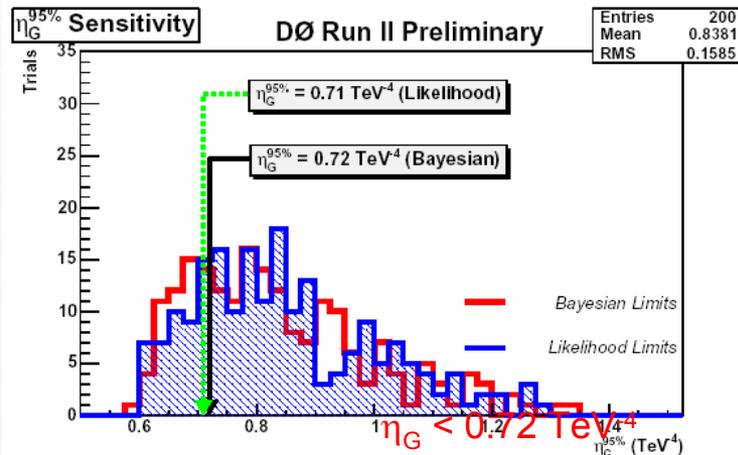


DØ Run II Preliminary

SM + ED terms ($\eta_G = 3.0 \text{ TeV}^{-4}$)



Source of systematics	Uncertainty
<i>K</i> -factor	10%
Choice of p.d.f.	5%
<i>pT</i> dependence on efficiency	5%
Choice of MC <i>pT</i> smearing	4%
MC to Data Normalization	1%
Total	13%



- Fit to 2D-templates via bilinear form in η_G
- First search in the dimuon channel at hadron colliders
- Current sensitivity approaches that for combination of all LEP channels
- Best limits on large ED in the dimuon channel have been set

Hewett		GRW	HLZ (TeV, @95% CL)					
$\lambda = +1$	$\lambda = -1$		$n = 2$	$n = 3$	$n = 4$	$n = 5$	$n = 6$	$n = 7$
0.97	0.95	1.09	1.00	1.29	1.09	0.98	0.91	0.86

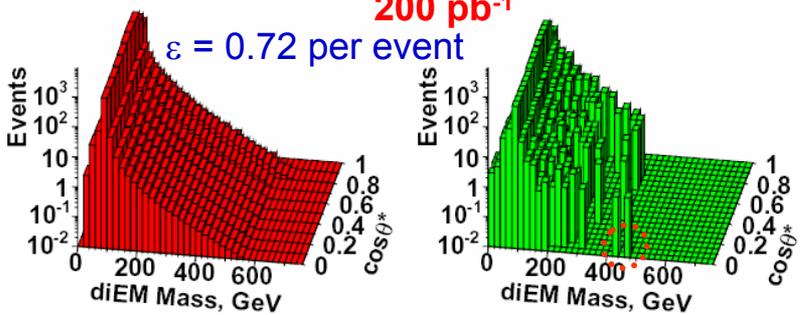


Search for Large ED in the diEM Channel

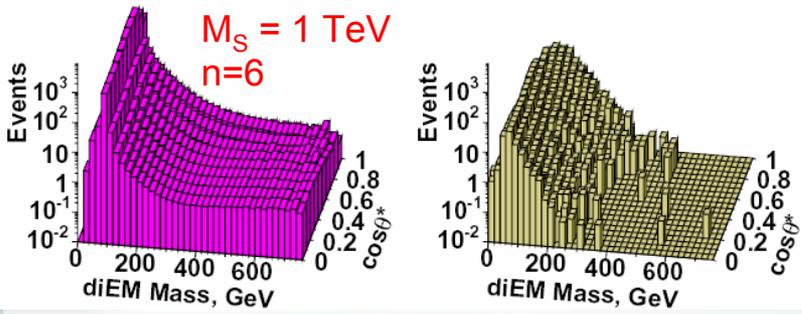
- Combine diphotons and dielectrons into “di-EM objects” to maximize efficiency
- High-mass, low $|\cos\theta^*|$ tail is a characteristic signature of large ED [Cheung & GL, PRD **62**, 076003 (2000)]

- Sensitivity is dominated by the diphoton channel ($2 \rightarrow 1 + 1$)
- Data agree well with the SM predictions; proceed with setting limits on large ED: alone or in combination with published Run I result [PRL **86**, 1156 (2001)]:

SM Prediction **DØ Run II Preliminary** **200 pb⁻¹** **Data**



ED Signal **QCD Background**



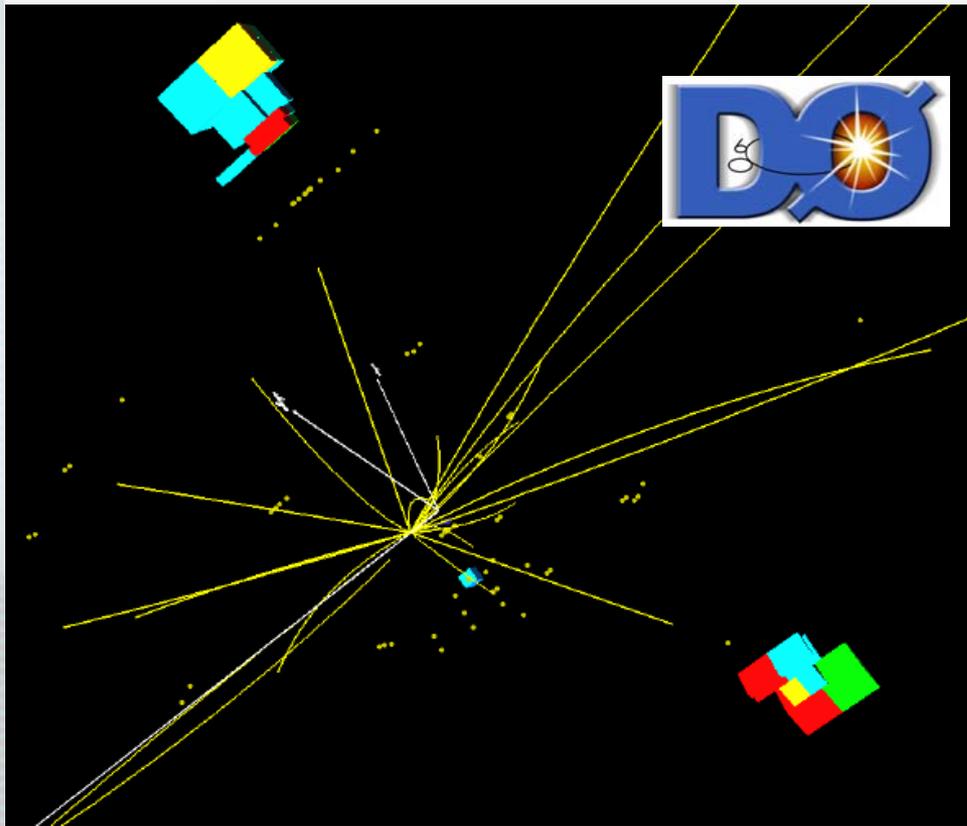
Hewett		GRW	HLZ (TeV, @95% CL)					
$\lambda = +1$	$\lambda = -1$		$n = 2$	$n = 3$	$n = 4$	$n = 5$	$n = 6$	$n = 7$
1.22	1.10	1.36	1.56	1.61	1.36	1.23	1.14	1.08
1.28	1.16	1.43	1.67	1.70	1.43	1.29	1.20	1.14
r_{\max}			170 μm	1.5 nm	5.7 pm	0.2 pm	21 fm	4.2 fm

- These are the most stringent constraints on large ED for $n > 2$ to date, among all the experiments
- For $n=2$, sensitivity is approaching that of the tabletop gravity measurements ($M_D = 1.7 \text{ TeV}$, $r < 160 \mu\text{m}$)

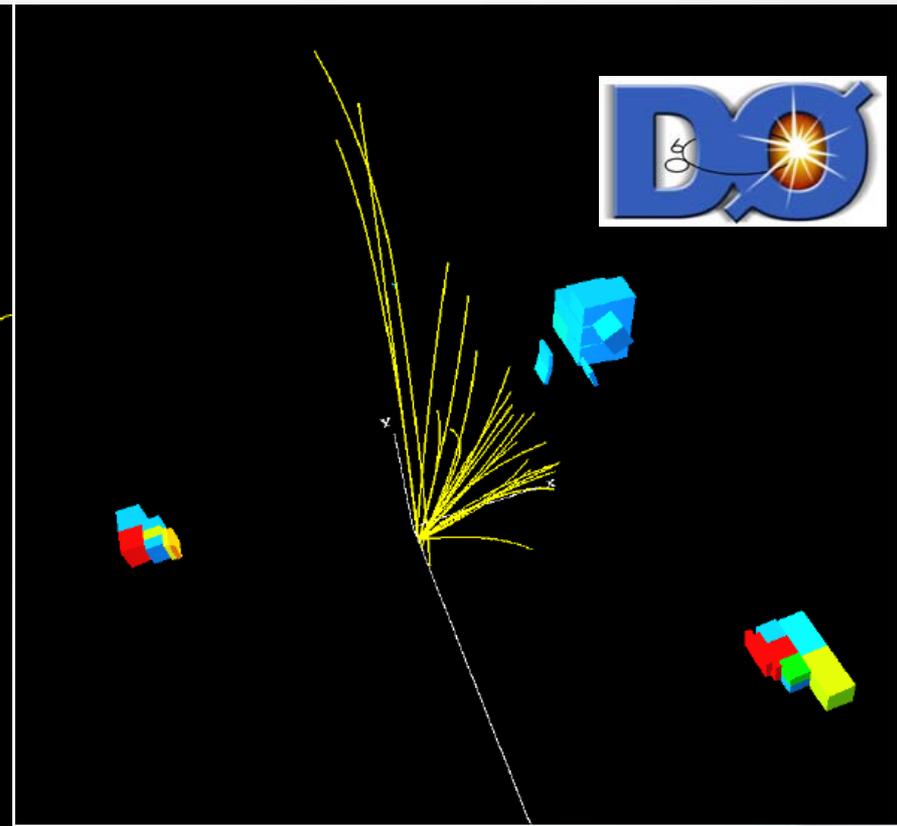


Interesting Candidate Events

- While the DØ data are consistent with the SM, the two highest-mass candidates have anomalously low value of $\cos\theta^*$ typical of ED signal:



Event Callas: $M_{ee} = 475 \text{ GeV}$, $\cos\eta^* = 0.01$



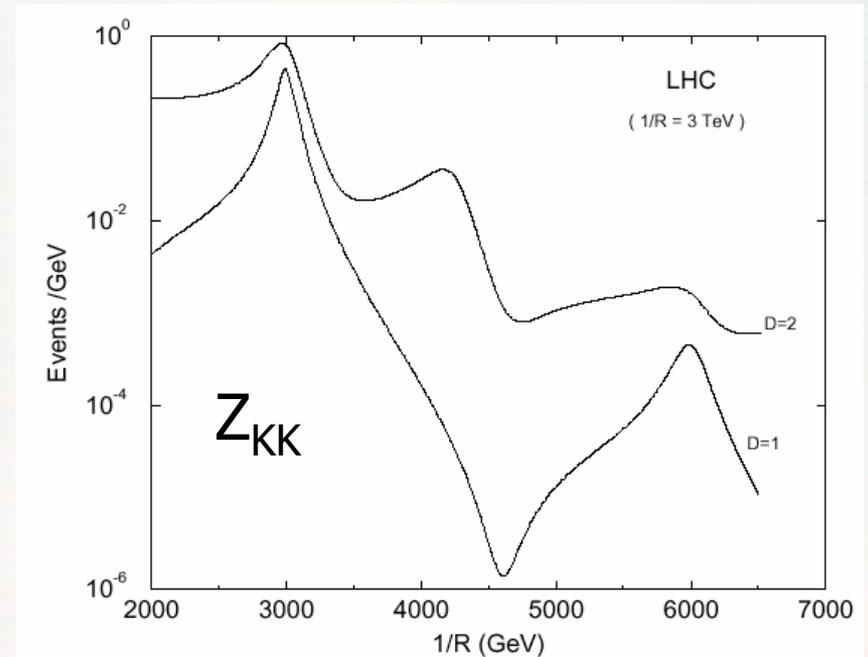
$M_{\gamma\gamma} = 436 \text{ GeV}$, $\cos\eta^* = 0.03$



TeV⁻¹ Extra Dimensions

- Intermediate-size **extra dimensions** with $\sim \text{TeV}^{-1}$ radius
- Introduced by [Antoniadis, PL **B246**, 377 (1990)] in the string theory context
- Used by [Dienes, Dudas, and Gherghetta, PL **B436**, 55 (1998)] to allow for low-energy unification
 - Expect Z_{KK} , W_{KK} , g_{KK} resonances at the LHC energies
 - At lower energies, can study effects of virtual exchange of the Kaluza-Klein modes of vector bosons
- Current indirect constraints come from precision EW measurements:
 $1/r \sim 6 \text{ TeV}$
- No dedicated experimental searches at colliders to date

[Antoniadis, Benaklis, and Quiros, PL **B460**, 176 (1999)]

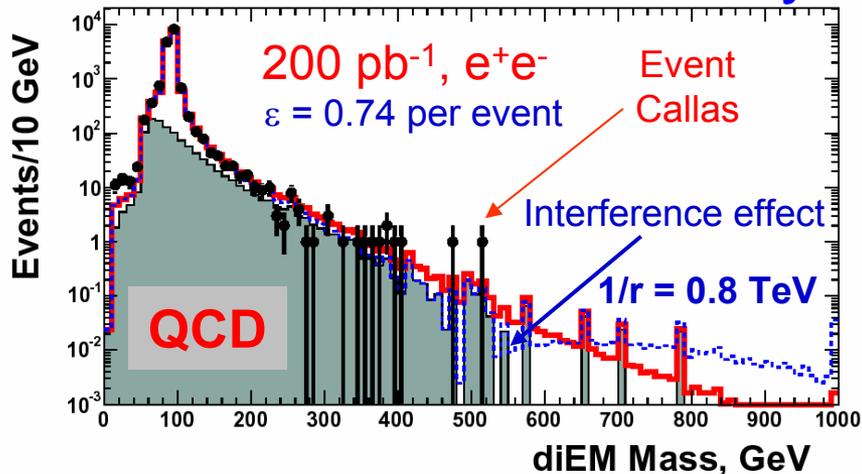




First Dedicated Search for TeV^{-1} ED

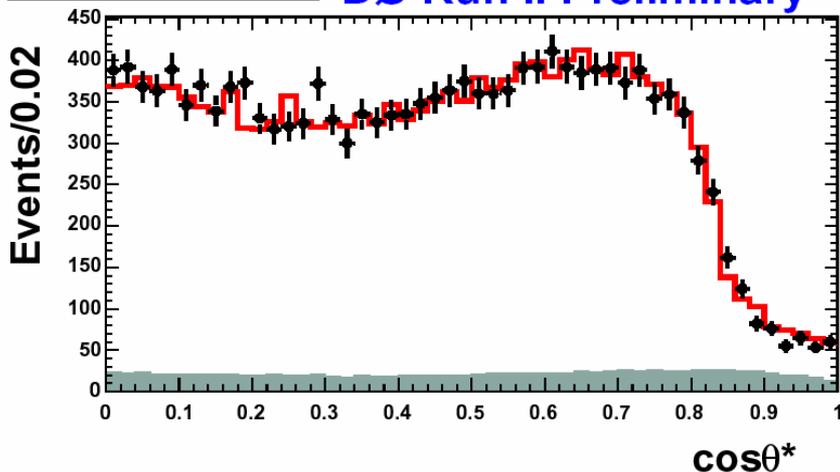
diEM Mass Spectrum

DØ Run II Preliminary



diEM $\cos\theta^*$ Spectrum

DØ Run II Preliminary



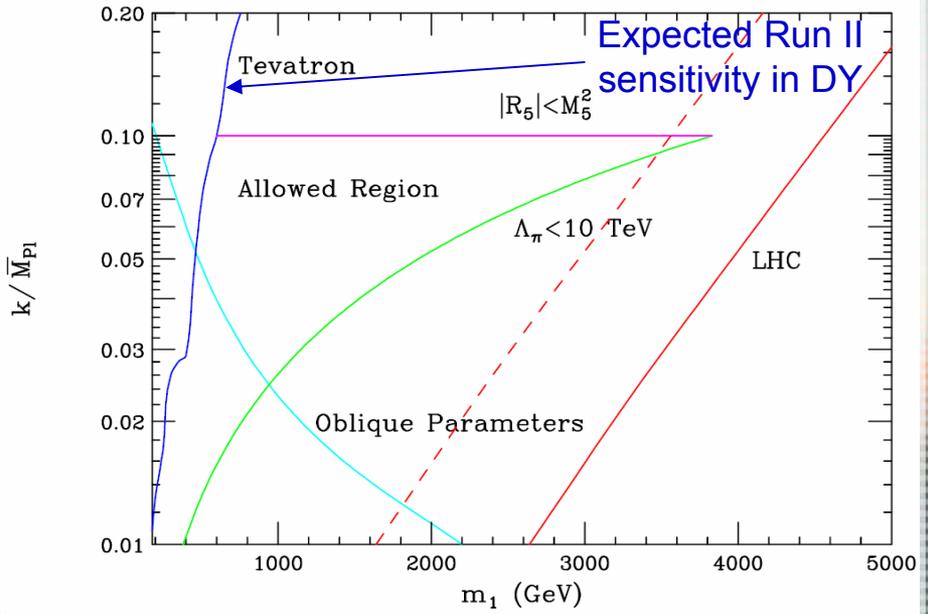
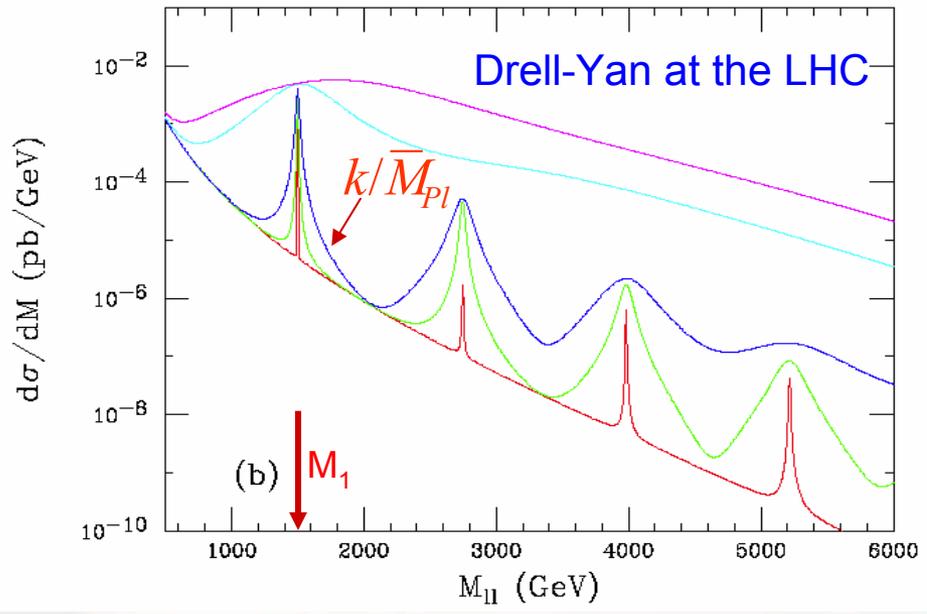
- While the Tevatron sensitivity is inferior to the indirect limits, it explores the effects of virtual KK modes at higher energies, i.e. complementary to those in the EW data
- DØ has performed the first dedicated search of this kind in the dielectron channel based on 200 pb^{-1} of Run II data ($Z_{\text{KK}}, \gamma_{\text{KK}} \rightarrow e^+e^-$)
- Based on the same diEM sample, with an additional requirement of at least one track
- The 2D-technique similar to the search for ADD effects in the virtual exchange yields the best sensitivity in the DY production [Cheung, GL PRD **65**, 076003 (2002)]
- Data agree with the SM predictions, which resulted in the following limit on their size:
 - ❖ $1/r > 1.12$ TeV @ 95% CL
 - ❖ $r < 1.75 \times 10^{-19}$ m



Randall-Sundrum Model Parameters

- Need only **two parameters** to define the model: k and r_c
- **Equivalent set** of parameters:
 - The mass of the first KK mode, M_1
 - Dimensionless coupling k / \bar{M}_{Pl}

- To avoid fine-tuning and non-perturbative regime, **coupling can't be too large or too small**
- $0.01 \leq k / \bar{M}_{Pl} \leq 0.10$ is the **expected range**
- **Gravitons are narrow**

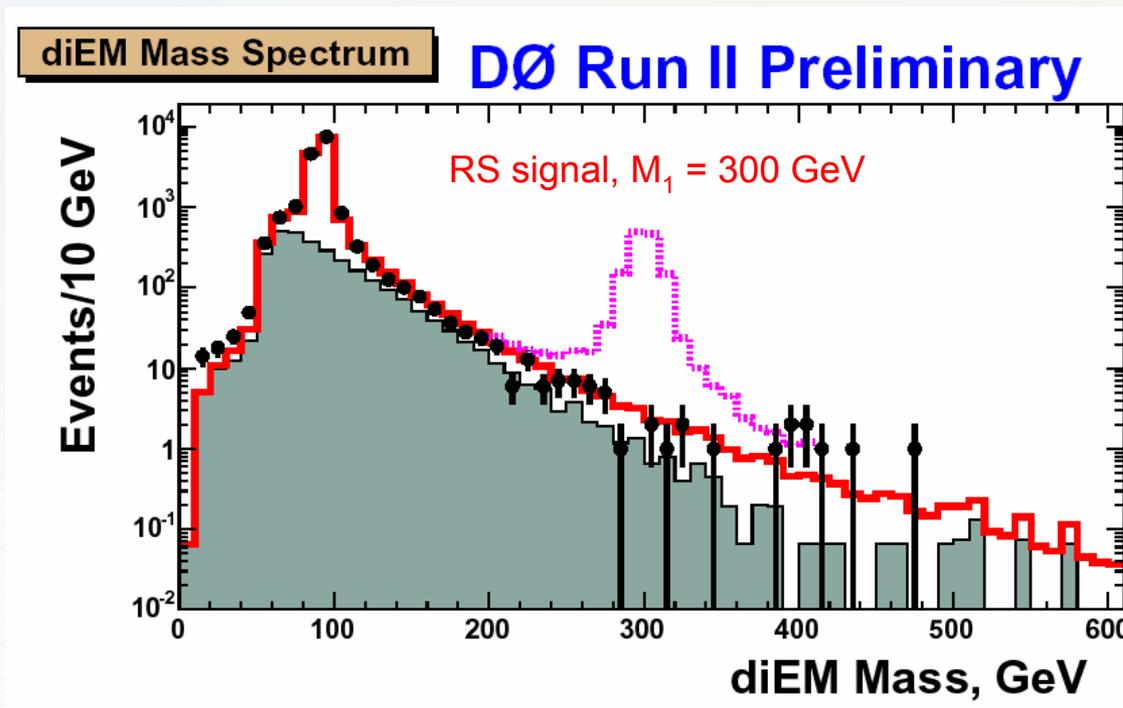


[Davoudiasl, Hewett, and Rizzo, PRD **63**, 075004 (2001)]



Search for RS Gravitons

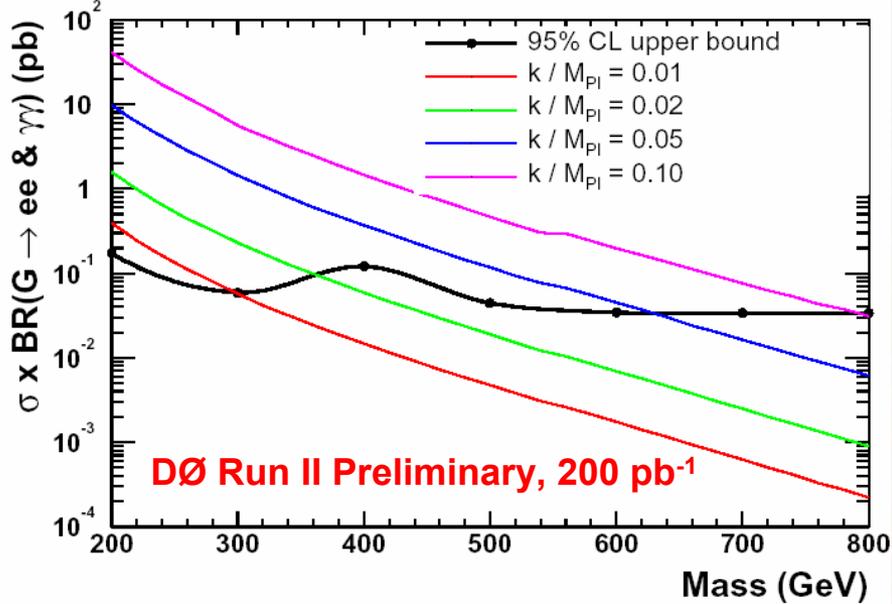
- **DØ** has just **completed preliminary analysis** and produced first results in August
- Analysis **based on 200 pb⁻¹ of e⁺e⁻ and $\gamma\gamma$ data** – the same data set as used for searches for large ED
- **“Naturally” combine two channels** by NOT distinguishing between electrons and photons; $\varepsilon = 0.85$ per EM object or 0.72 per event!
- **Search window size has been optimized** to yield maximum signal significance





DØ Limits in the $ee+\gamma\gamma$ Channel

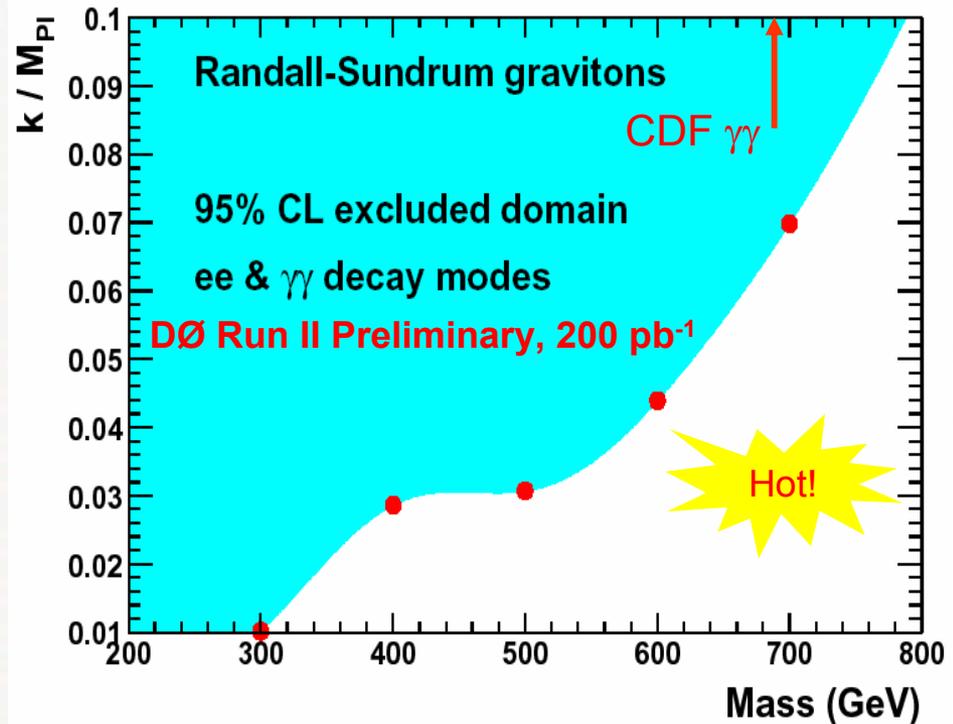
The tightest limits on RS gravitons to date



Already better limits than the sensitivity for Run II, as predicted by theorists!

Assume fixed K -factor of 1.3 for the signal

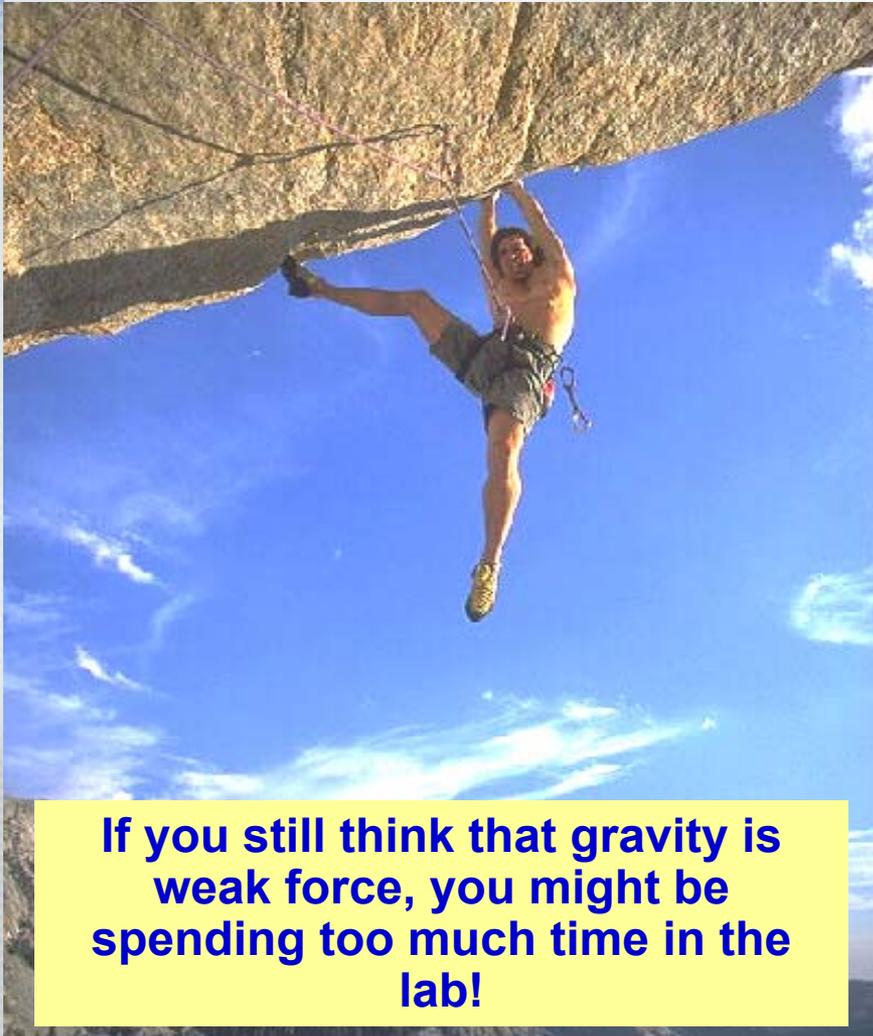
Masses up to 780 GeV are excluded for $k/M_{Pl} = 0.1$



- While the data agree well with the SM, a few interesting candidate events have been observed
 - N.B.: The three highest-mass events: ee , $\mu\mu$, and $\gamma\gamma$ – all have low value of $|\cos\theta^*|$!
- We expect an increase in sensitivity to ED models by a factor of ~ 2 within the next few years of the Tevatron operations, which would bring us into interesting area of the allowed parameter space
- Twice the analyzed statistics is already on tape, and more will come soon
- Then it'd be up to the LHC to either disprove or confirm the existence of extra dimensions in space with low-energy phenomenology



Conclusions



If you still think that gravity is weak force, you might be spending too much time in the lab!

- Colliders offer **ultimate probe** of models with **extra spatial dimensions**
- **Tevatron** is the highest-energy machine today and **has an excellent opportunity** to find first signs of **extra dimensions** in space
- **Accelerator performance is excellent** and both **CDF and DØ** have a large number of **new results** in this area
- Sensitivities **beyond the existing limits** have been achieved already:
 - **Tightest limits on large ED to date**
 - **Best limits on RS gravitons**
 - **Quick progress towards monojet searches with superior sensitivity**
- **Stay tuned** for more exciting results to come with larger data sets – any day now could bring a **dramatic discovery!**



P.S. Fine Tuning Explained...

• Fine tuning explained:

– Numerology: $987654321/123456789 =$

$8.000000073 ?$

• Numerology it is not!

$$\lim_{N \rightarrow \infty} \frac{NML\dots987654321}{123456789\dots LMN} = N - 2$$

– Seeing is believing:

• In hexadecimal system,

$FEDCBA987654321/123456789ABCDEF =$

14.00000000000000000183