

Energy Conservation and Pomeron Loops in High Energy Evolution

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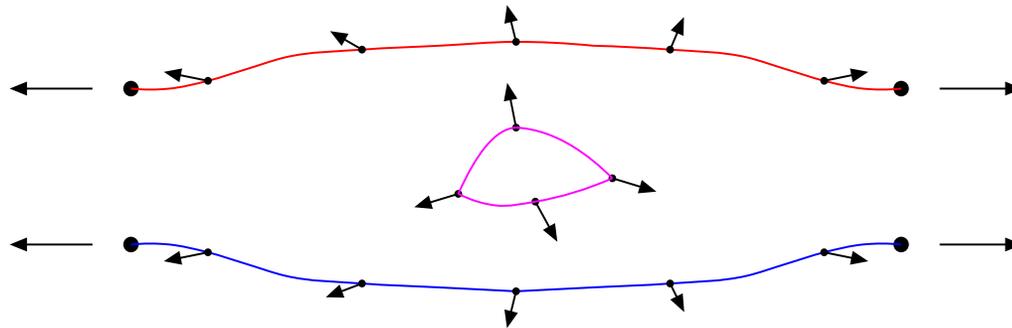
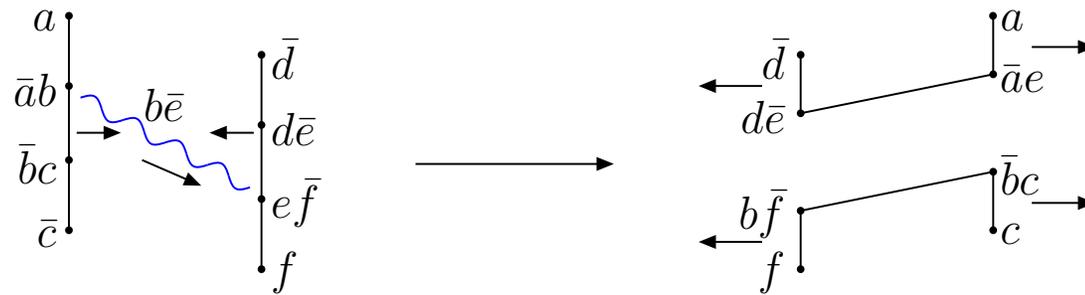
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Outline

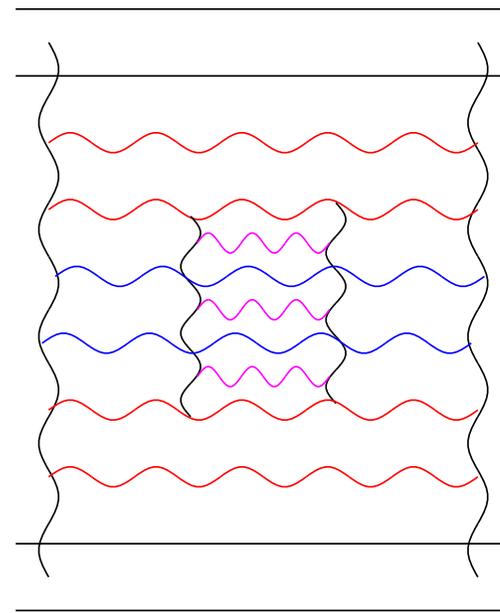
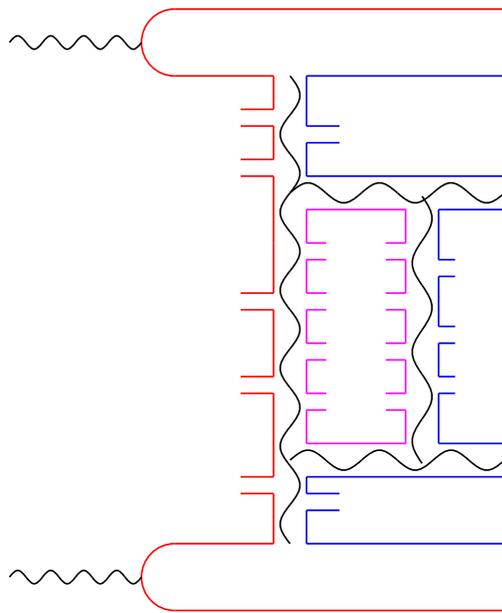
- Dipoles in High Energy QCD.
- Improving Mueller's Dipole Model:
 - Energy-Momentum Conservation.
 - Colour Suppressed Effects.
- DIS at HERA.
- pp at the Tevatron.
- Summary and Outlook.

Strings and Dipoles in HE Collisions



From Dipoles to Pomeron

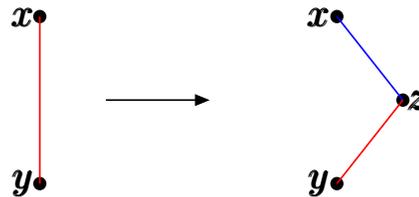
- Multiple interactions will generate colour loops. The colour loops in the inelastic amplitude, when squared, give rise to pomeron loops in the elastic amplitude.



Mueller's Dipole Model

- In the dipole model the splitting kernel is given by

$$\mathcal{M}(\mathbf{x}, \mathbf{y}, \mathbf{z}) = \frac{\bar{\alpha}_s}{2\pi} \frac{(\mathbf{x} - \mathbf{y})^2}{(\mathbf{x} - \mathbf{z})^2 (\mathbf{z} - \mathbf{y})^2}$$



- This expression diverges when integrated over d^2z . A cutoff, $\rho^2 \leq (\mathbf{x} - \mathbf{z})^2, (\mathbf{y} - \mathbf{z})^2$, is needed to regulate the divergence.
- Even though the divergence cancels against virtual emissions, the cut off, ρ , has to be kept in a MC.

Conserving Energy-Momentum

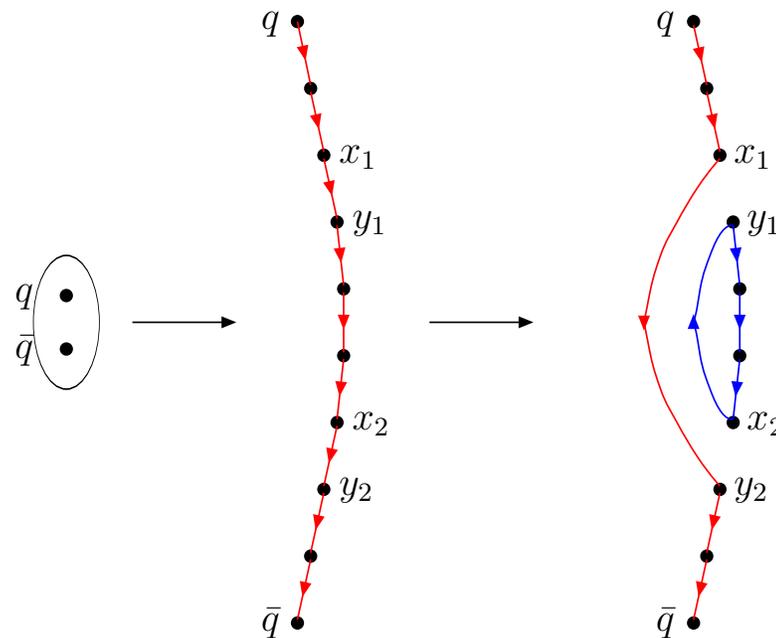
- A small dipole corresponds to two well localized gluons.
⇒ large $p_T \sim 1/r$ for these gluons. The emission of such gluons violates EM conservation.
- These emissions are compensated by virtual emissions.
⇒ σ_{tot} is determined by real emissions. Keeping only these we get a closer correspondence between the cascade and exclusive final states.
- Effects of EM conservation: hep-ph/0503181, JHEP 2005.

Beyond Large N_c

- Beyond large N_c , higher order multipoles are formed, such as quadrupoles.
- Try to approximate a quadrupole by two dipoles
 - Necessary to include a “dipole swing”.
- Multiple Interactions suppressed by N_c^2 while evolution leading in N_c
 - ⇒ Formalism not frame independent.

Dipole Swing

- Include a $2 \rightarrow 2$ transition (suppressed by N_c^2). This can be seen as a dipole swing (or a colour recoupling): $(x_1, y_1), (x_2, y_2) \rightarrow (x_1, y_2), (x_2, y_1)$.



- Result essentially frame independent.

The γ^*p Cross Section

- To simulate DIS events we picture the γ^* fluctuate into a $q\bar{q}$ pair which initiates the dipole cascade.
- The $\gamma^* \rightarrow q\bar{q}$ vertex is given by the wavefunctions $\psi_T(z, \mathbf{r})$ and $\psi_L(z, \mathbf{r})$.

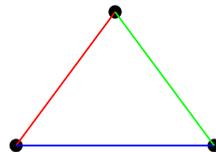
$$\sigma_{\gamma^*p}^{tot} = \sigma_T + \sigma_L = \int d^2\mathbf{r} \int_0^1 dz (|\psi_T|^2 + |\psi_L|^2) \sigma(z, \mathbf{r})$$

$$\sigma(z, \mathbf{r}) = 2 \int d^2\mathbf{b} \langle 1 - \exp(-\sum_{ij} \mathcal{T}_{ij}) \rangle$$

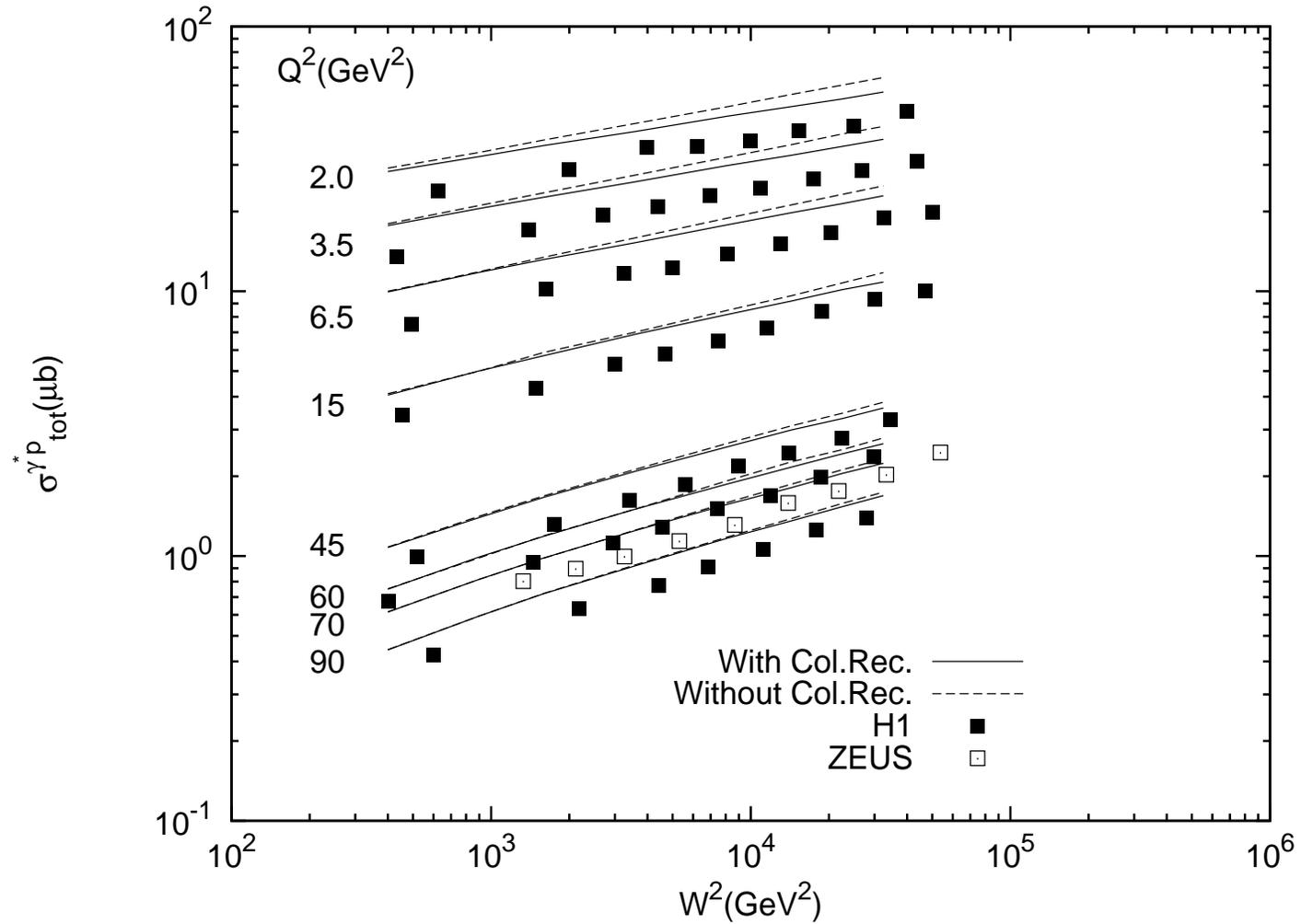
$$\mathcal{T}_{ij} = \frac{\alpha_s^2}{8} \log^2 \left\{ \frac{(\mathbf{x}_i - \mathbf{y}_j)^2 (\mathbf{y}_i - \mathbf{x}_j)^2}{(\mathbf{x}_i - \mathbf{x}_j)^2 (\mathbf{y}_i - \mathbf{y}_j)^2} \right\}$$

The Proton

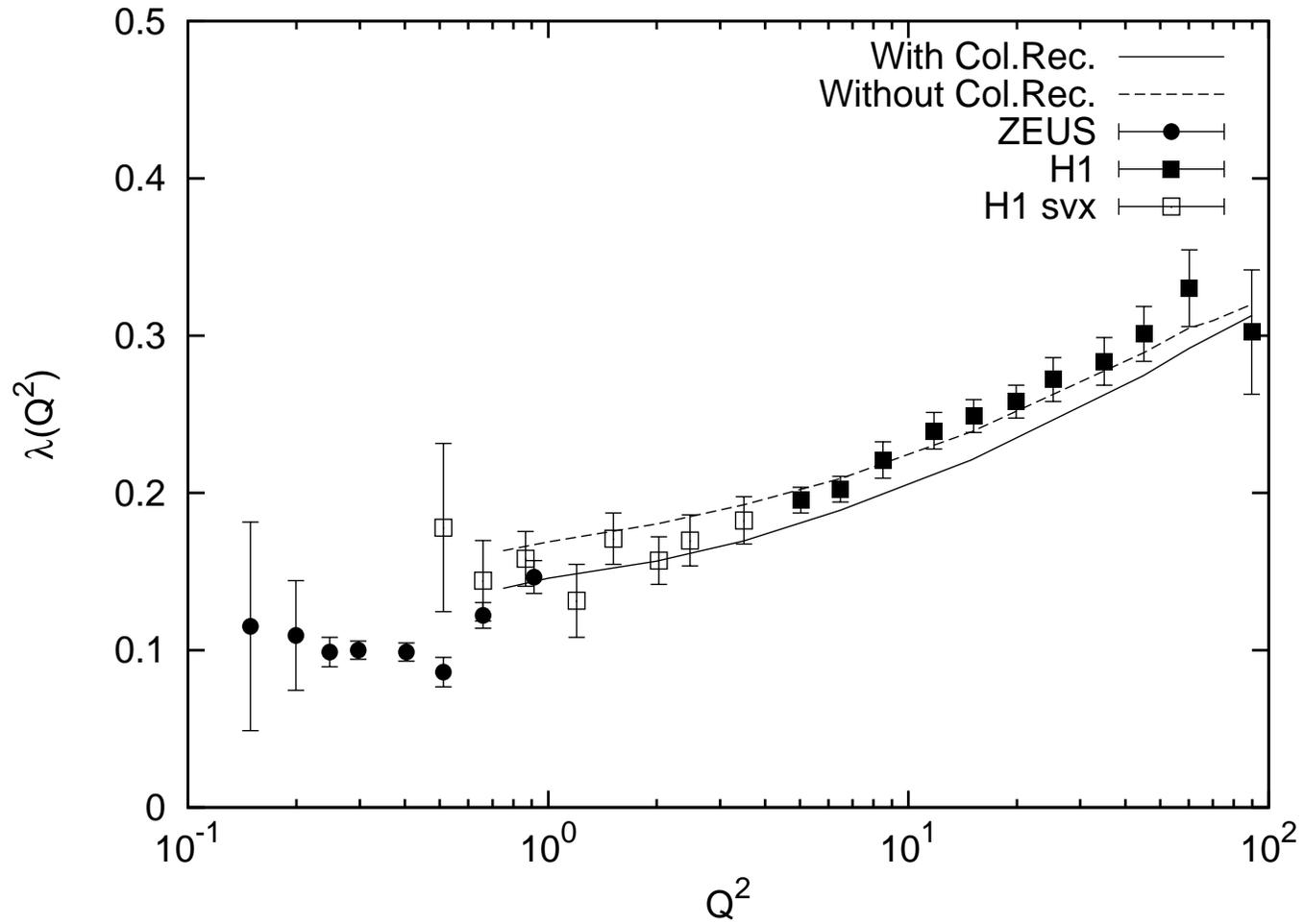
- Soft physics involved, cannot be fully described by PQCD.
- High energy pp collision \implies scattering dominated by minijets = hard subcollisions.
- For the first gluon emission the proton appears as 3 dipoles.
 - However, no factorization, even in large N_c limit.
- As an approximation we use a Δ shaped topology consisting of 3 colour dipoles.



Results for DIS

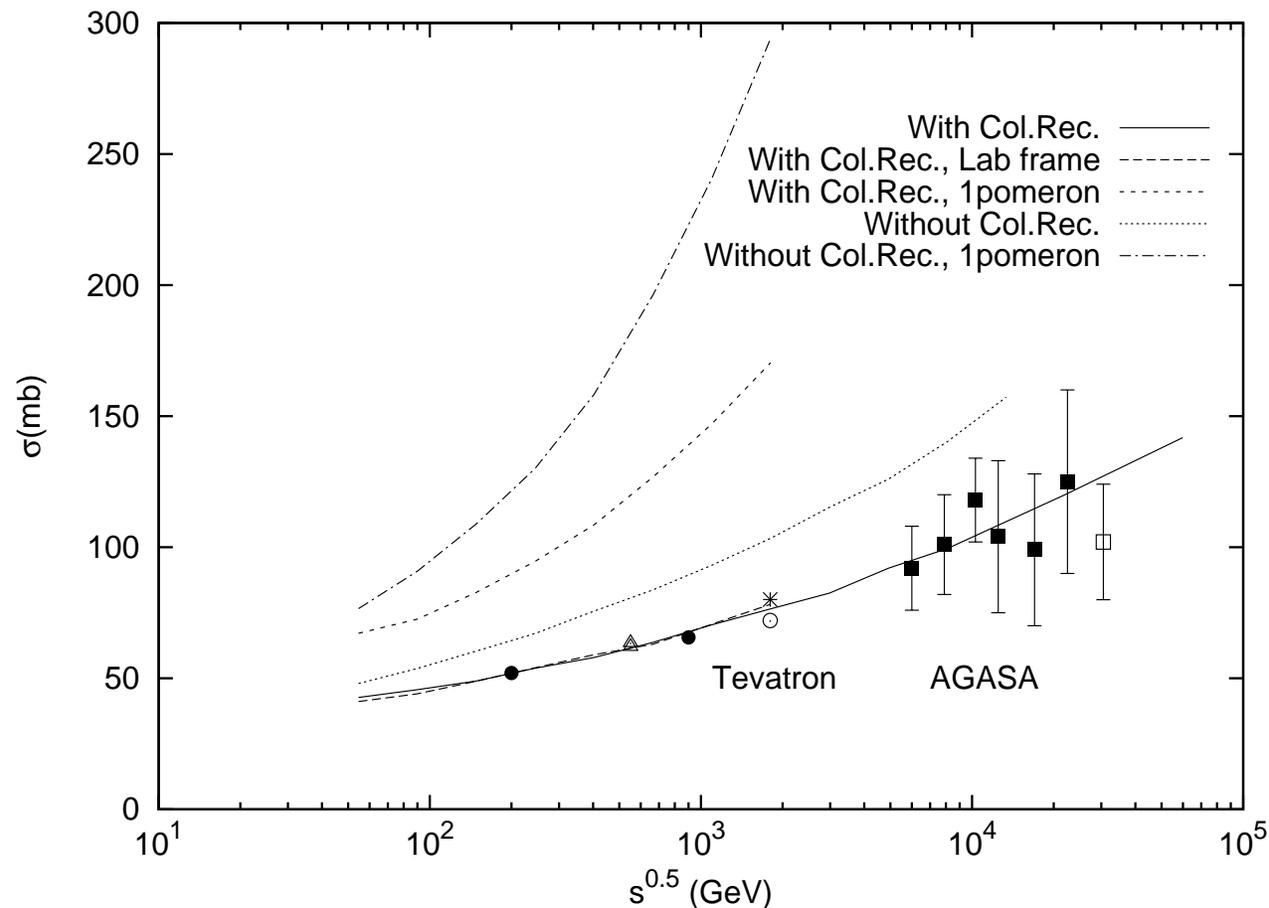


F_2 slope at HERA

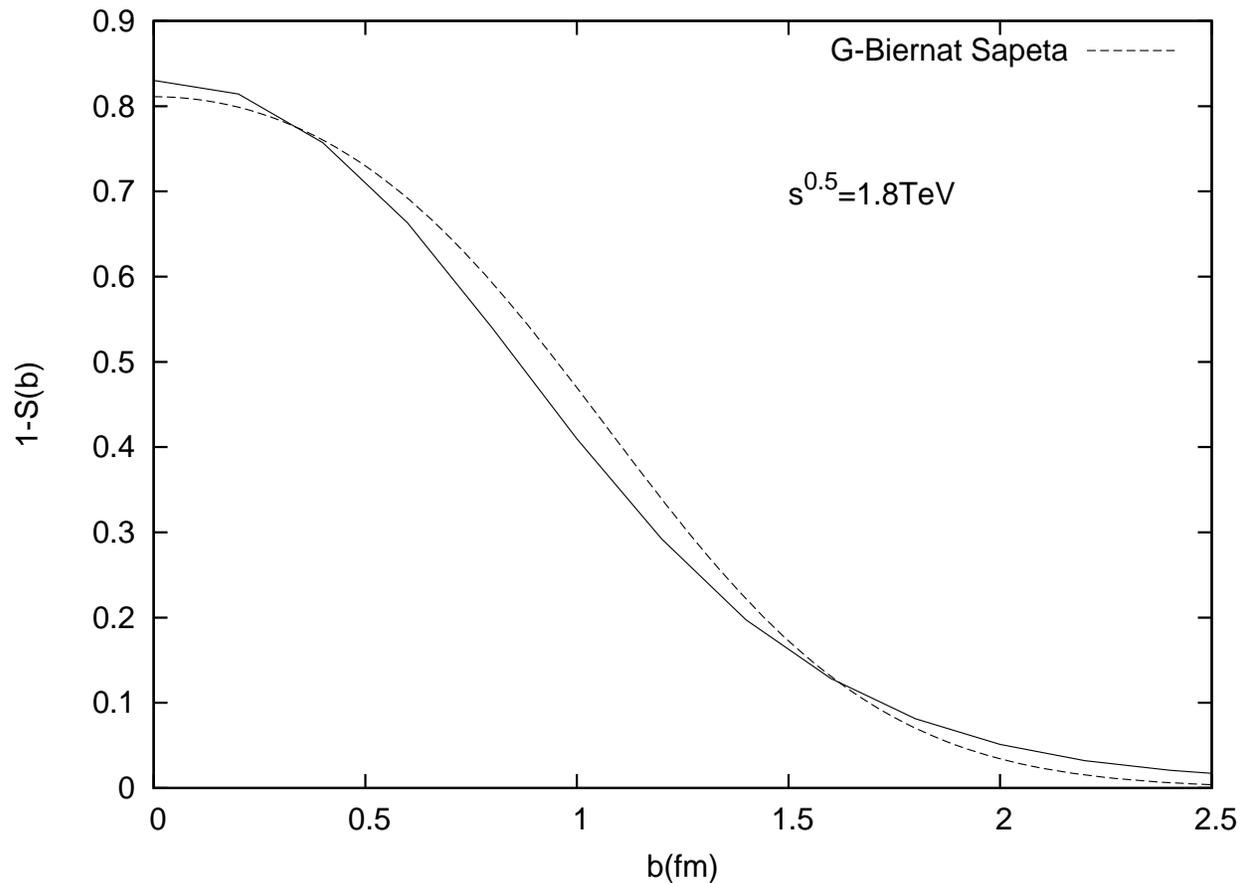


pp Total Cross Section

Saturation very important (small effect at HERA).
Dipole swing \implies Near frame independence.



pp Impact Parameter Profile



$S = \exp(-\sum_{ij} \mathcal{T}_{ij})$. GB-S = two parameter fit to data.

Summary and Outlook

- Including EM conservation in the Dipole Model we get a good agreement with data for DIS at HERA.
- We have included some effects beyond leading N_c
 - These are seen to be very important for pp collisions.
 - Gives good agreement with data up to Tevatron energies.
- Obtain an explicitly frame independent formalism for subleading N_c effects in dipole language.
- Obtain results for exclusive final states.